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Ultrahigh-Q microcavity photonics and optomechanical cooling

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清华大学高等研究院 2015年3月11日

Microcavity photonics and quantum optics group@PKU
URL: www.phy.pku.edu.cn/~yfxiao/

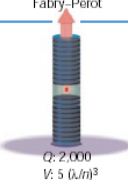
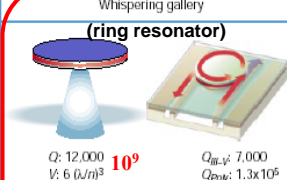
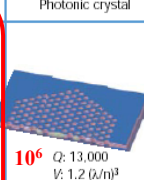

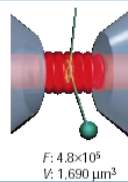
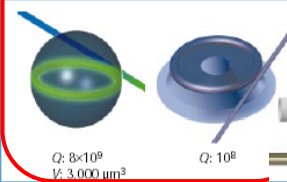
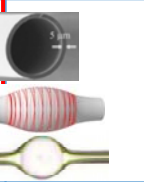

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Light-matter interaction plays the central role in both fundamental optical physics and applied photonics, ranging from quantum light source to functional photonic device.

2

Optical microresonator – important platform for... Confinement

Photon Lifetime	High Q	 Fabry-Perot $Q: 2,000$ $V: 5 (\mu\text{m})^3$	 Whispering gallery (ring resonator) $Q: 12,000$ 10^9 $V: 6 (\mu\text{m})^3$ $Q_{\text{in}}: 7,000$ $Q_{\text{out}}: 1.3 \times 10^5$	 Photonic crystal 10^6 $Q: 13,000$ $V: 1.2 (\mu\text{m})^3$	 Plasmonic microcavity 500 nm
	Ultrahigh Q	 $F: 4.8 \times 10^5$ $V: 1.690 \mu\text{m}^3$	 $Q: 8 \times 10^8$ $V: 3,000 \mu\text{m}^3$ $Q: 10^8$	 	 Xiao et al., PRL, 2010


K. J. Vahala Nature 424, 839

1. FP-type microcavity
2. Whispering gallery microcavity
3. Photonic crystal microcavity
4. Plasmonic microcavity


Ultrahigh Q, very small V
mass production on a chip

3

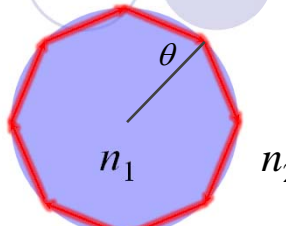
Whispering-gallery microcavities




St. Paul's Church




Temple of Heaven

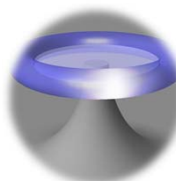




Whispering Gallery




Echo Wall





acoustic wave → micro wave → optical wave

4

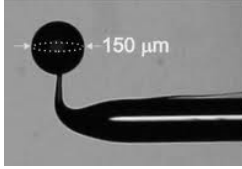

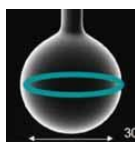

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Optical WGMs in liquid and silica microspheres

Early study: liquid droplets


From liquid to solid state

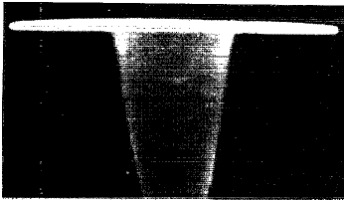
many phenomena, e.g., nonlinear optical process but, suffered from evaporation

- C. G. B. Garrett, W. Kaiser, and W. L. Bond, *Phys. Rev.* 124, 1807 (1961).
- A. Ashkin and J. M. Dziedzic, *Phys. Rev. Lett.* 38, 1351 (1977).
- P. Chylek, J. T. Kiehl, and M. K. W. Ko, *Phys. Rev. A* 18, 2229 (1978).
- V. B. Braginsky, M. L. Gorodetsky, and V. S. Ilchenko, *Phys. Lett. A* 137, 393 (1989).
- M. C. Kuzyk and H. Wang, *SPIE Newsroom*. DOI: 10.1117/2.1201208.004434

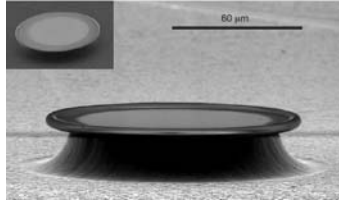
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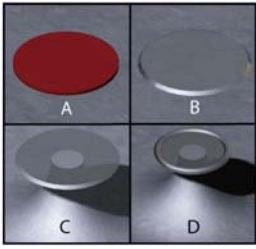
WGMs in silica microdisk and microtoroid on a chip



S. L. McCall et al., *Appl. Phys. Lett.* 60, 289 (1992)




Vahala Group, *Nature* 421, 925-927 (2003)

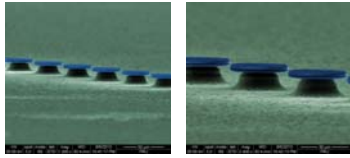


- **Microtoroid resonator fabrication**
 - Photolithographically defined photo-resist disks
 - Silica etching using buffered HF and PR removal leaving silica disk
 - XeF₂ isotropical etch of Si and silica disk undercut
 - CO₂ laser activated reflow to form microtoroid

6

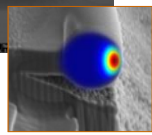

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Why ultra-high-Q microcavities on a silicon chip



@Peking

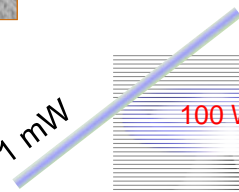
Diameter: 20 μm – 200 μm
 Quality factor > 10^8 ,
 Intracavity lifetime > 100 ns

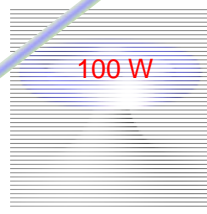


“Strongly enhanced light-matter interaction”

$P_{\text{in}} = 1 \text{ mW} \rightarrow$
 $P_{\text{cav}} \sim 100 \text{ W}, I_{\text{cav}} > 2 \text{ GW/cm}^2,$
 $\tau \sim 100 \text{ ns} \rightarrow \# \text{ of round trip} > 10^5.$


1 mW



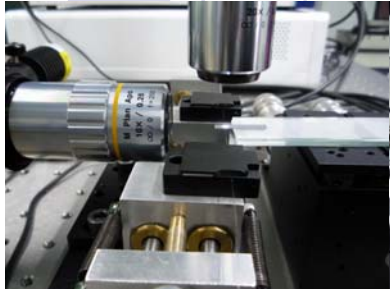


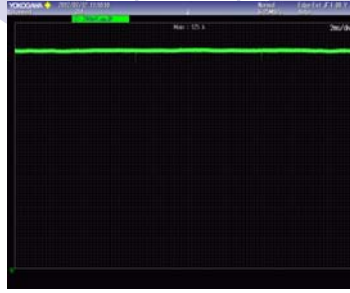
100 W

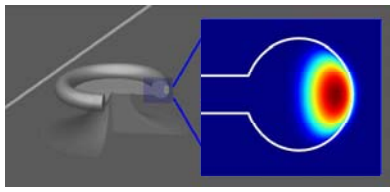
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Our Group






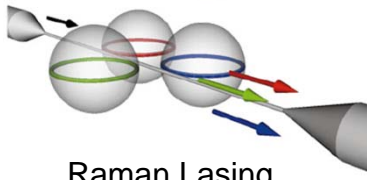


Linewidth < 2 MHz
Quality factor > 2×10^8
From under coupling, critical coupling, to over coupling

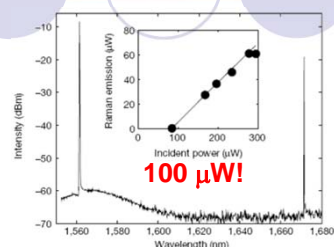
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APPLICATION: Low-threshold lasing



Raman Lasing



Intensity (dBm)

Raman emission (μW)

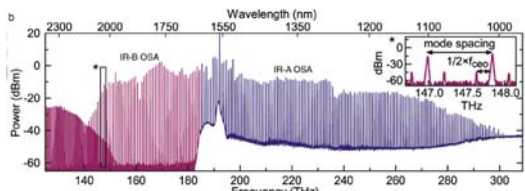
Incident power (μW)

100 μW!

Wavelength (nm)

Vahala group, Nature (2002)

Single-frequency Pump → Microcomb



Power (dBm)

Frequency (THz)

Wavelength (nm)

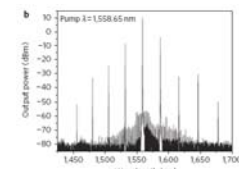
IR-B OSA

IR-A OSA

mode spacing

$1/2 \times f_{\text{cav}}$

147.0 147.5 148.0 THz



Output power (dBm)

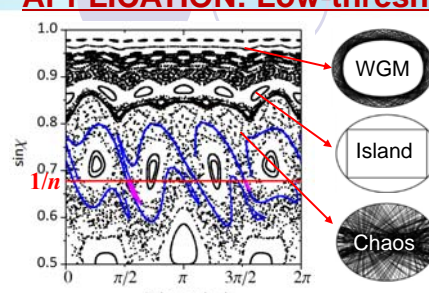
Wavelength (nm)

Pump $\lambda = 1558.65$ nm

Kippenberg group, PRL (2008)
L. Razzari *et al.*, Nature Photon. 4, 41 (2010)

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APPLICATION: Low-threshold lasing



WGM


Island

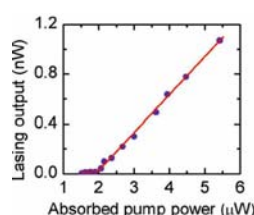
Chaos

$1/n$

Polar angle ϕ

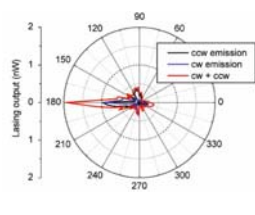
Question: Ultrahigh-Q factor
+
Highly directional emission





Lasing output (nW)

Absorbed pump power (μW)



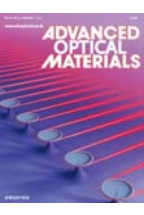
Lasing output (nW)

Polar angle (degrees)


— cw emission

— cw emission

— cw + cw



ADVANCED OPTICAL MATERIALS




Optik & Photonik

Opt. & Photon. News

X. F. Jiang, Y.-F. Xiao* *et al.*, *Advanced Materials* 24, OP260-OP264 (2012)

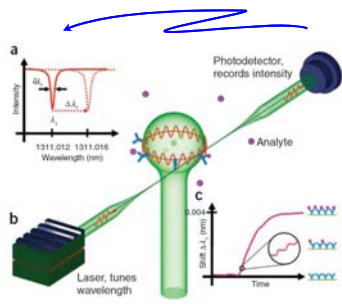
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APPLICATION: Highly sensitive optical biosensing

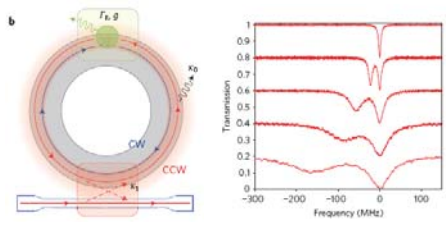
Mode shift



Thermal and probe noise

Mode splitting

Rayleigh scattering
 → mode coupling
 → mode splitting (splitting > linewidth)
 (immune to thermal noise)



F. Vollmer & S. Arnold, Nature Methods 5, 591 (2008)
 M. A. Santiago-Cordoba et al., APL 99, 073701 (2011)
 J. D. Swaim, J. Knittel, and W. P. Bowen, APL 99, 243109 (2011)
 V. R. Dantham et al., APL 101, 043704 (2012)
 Y.-F. Xiao et al., Phys. Rev. A 85, 031805(R) (2012)
 V. R. Dantham et al., Nano Lett. 13, 3347 (2013)



J. Zhu et al. Nature Photonics 4, 46 (2010).
 L. He et al., Nature Nanotech. 6, 428 (2011).
 A. Mazzei et al. Phys. Rev. Lett. 99, 173603 (2007).
 L. Chantada et al., JOSA B 25, 1312 (2008).
 X. Yi, Y.-F. Xiao* et al., APL 97, 203705 (2010); PRA 83, 023802 (2011); JAP 111, 114702 (2012)

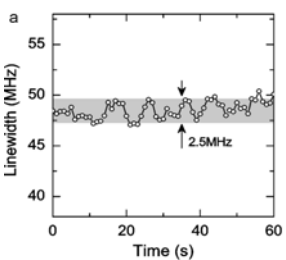
11

APPLICATION: Highly sensitive optical biosensing

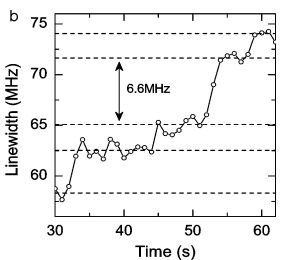
Mode broadening:

- Temperature insensitive
- High sensitivity
- No requirement of ultra-high Q factors







Pure water (w/o virus)




PS nanoparticles



L. Shao, X.-F. Jiang (Y.-F. Xiao*) et al., Advanced Materials 25(39), 5615 (2013)

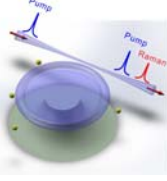
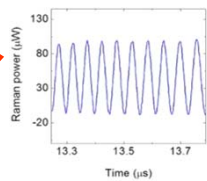
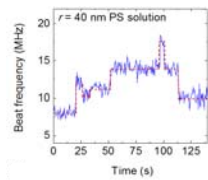
12

APPLICATION: Highly sensitive optical biosensing

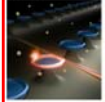


Microcavity Raman laser:

- Temperature insensitive
- Higher sensitivity than passive cavity
- No requirement of doping

Physicists develop miniature Raman laser sensors for single nanoparticle detection




[The rest of the article text is too small to transcribe accurately]

B.-B. Li, Y.-F. Xiao* et al., PNAS 111, 14657 (2014)

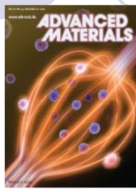
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APPLICATION: Highly sensitive optical biosensing

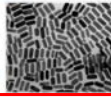


Nanofiber array sensor:

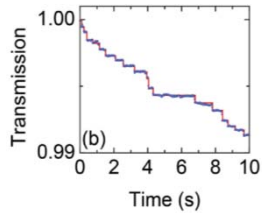
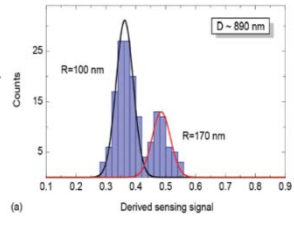
- No requirement of tunable laser
- Single nanoparticle response
- Larger sensor area



Nanofiber sensor can detect and size a single nanoparticle



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Innovative Nanofiber-Array Sensor Detects and Sizes Individual Nanoparticles

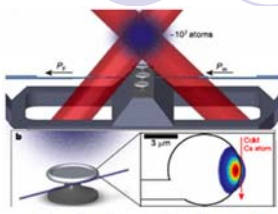
[The rest of the article text is too small to transcribe accurately]

X.-C. Yu, Y.-F. Xiao* et al., Adv. Mater. 26, 7462 (2014)

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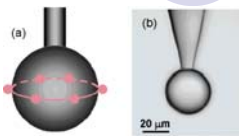
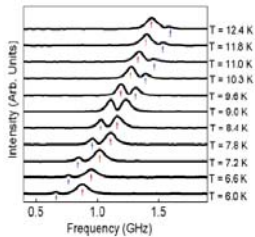
APPLICATION: Quantum optics & quantum information



Strong Coupling on a Microelectronic Chip

“Strong coupling”

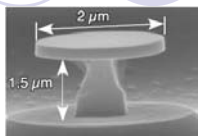
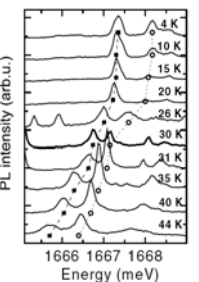
$g/2\pi \sim 50$ MHz
 $\kappa/2\pi \sim 18$ MHz
 $\gamma/2\pi \sim 2.6$ MHz

Intensity (Arb. Units)

Frequency (GHz)

$T = 12.4$ K
 $T = 11.8$ K
 $T = 11.0$ K
 $T = 10.3$ K
 $T = 9.6$ K
 $T = 9.0$ K
 $T = 8.4$ K
 $T = 7.8$ K
 $T = 7.2$ K
 $T = 6.6$ K
 $T = 6.0$ K

PL Intensity (arb.u.)

Energy (meV)

4 K
 10 K
 15 K
 20 K
 26 K
 30 K
 31 K
 35 K
 40 K
 44 K

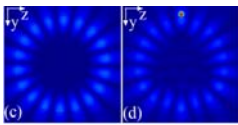
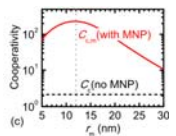
Cavity QED for implementation of quantum computation, quantum routers, ...

Nature **443**, 671 (2006); *Science* **319**, 1062 (2008); *PRL* **95**, 067401 (2005); *Nano Lett.* **6**, 2075 (2006)

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APPLICATION: Quantum optics & quantum information

Plasmonics enhanced interaction

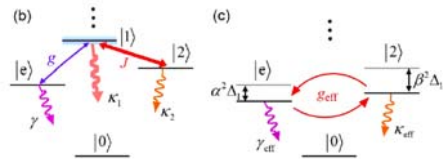
Cooperativity

r_c (nm)

C_{co} (with MNP)


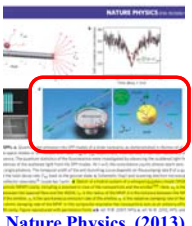
C (no MNP)

Strong coupling in coupled dissipative cavities



(b) (c)

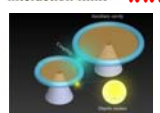
$|e\rangle$, $|1\rangle$, $|2\rangle$, $|0\rangle$
 κ_1 , κ_2 , γ , $\alpha^2 \Delta$, $\beta^2 \Delta$, γ_{eff} , κ_{eff}

Nature Physics (2013)

Jun 10, 2014

New design surpasses the coherent interaction limit www.phys.org




PHYS.ORG

Coherent light-matter interaction of the single-photon level is a long-standing goal for quantum operations. To reach such strong coupling regime, previous studies mostly focus on improving the quality (Q) factor and reducing the mode volume (V) of the same optical resonators, but they have already reached their limits due to current fabrication and other technological restrictions. Very recently, a team led by Professor Xiao Liu and Long-Guang Liang of Peking University, collaborated with Christoph Langens, have demonstrated the strong coupling via dual-mode resonators in a coupled-cavity quantum electrodynamics system. By increasing the resonator for both Q and small V as

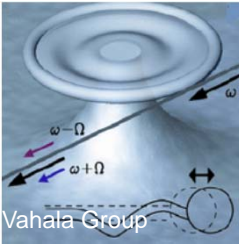
Xiao*, Liu* et al., Phys. Rev. A **85**, 031805(R) (2012) (Rapid Communications)

Liu et al., Phys. Rev. Lett. **112**(21), 213602 (2014)


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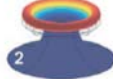
APPLICATION: Cavity optomechanics



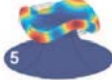
Vahala Group



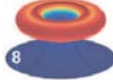
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2




5



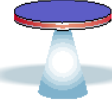
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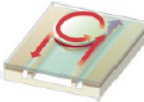

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Applications of whispering gallery microcavities

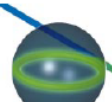
Whispering gallery




$Q: 12,000$
 $V: 6 (\lambda/n)^3$



$Q_{th}: 7,000$
 $Q_{poly}: 1.3 \times 10^5$



$Q: 8 \times 10^8$
 $V: 3,000 \mu m^3$



$Q: 10^8$

Ultrahigh Q, very small V,
 and mass production on a chip

Excellent platform for


Fundamental physics:

- Quantum optics
- Cavity quantum electrodynamics (QED)
- Quantum information
- Quantum chaos
- Cavity quantum optomechanics
-

Microphotonics:

- Highly sensitive bio/chemical sensing
- All-optical low-threshold switching
- On-chip microlasing (includ. Raman)
- Optical microcomb
- Optical filtering...


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Outline

1. Introduction
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Cavity optomechanics




Mechanical effect of light

Astronomic scale

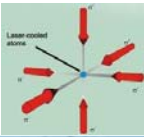
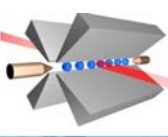



Comet's tail
Kepler

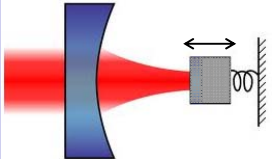
Microscopic scale

Laser cooling, Nobel 1997


Macro and Mesoscopic scale



Opto-mechanical coupled system

- Optical force acts on the mirror;
- Mirror's motion state changes;
- Cavity length changes;
- Cavity field changes;


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Applications

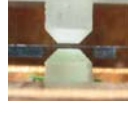
- Fundamental studies of quantum theory
 - Macroscopic quantum state, superposition and entanglement
 - Quantum-classical boundary
- Precision measurement
 - Gravitational waves detection
 - Mass, force and displacement
- Quantum information processing platform

MASS
 kg
 pg



Hz
 Mechanical frequency
 MHz


Kippenberg & Vahala
Science, 321, 1172 (2008)



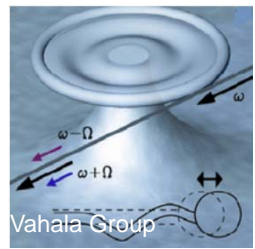
Trapped particle
Cold atoms/BEC

Weiping Zhang group,
PRL 108, 240405 (2012)

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WGM Cavity optomechanics



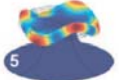
Vahala Group



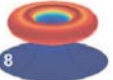
1



2




5



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Optomechanical coupling (dispersive)

C. K. Law et al., 1990s

System Hamiltonian $H = \omega_c(x)a^\dagger a + \omega_m b^\dagger b$

The cavity resonance frequency is modulated by the mechanical motion

$$\omega_c(x) = \omega_c + x \partial \omega_c(x) / \partial x + \mathcal{O}(x)$$

$x = x_{\text{ZPF}}(b^\dagger + b)$ **Zero-point fluctuation:** $x_{\text{ZPF}} = \sqrt{\hbar / (2m_{\text{eff}}\omega_m)}$

$$H_{\text{int}} = g a^\dagger a (b^\dagger + b)$$

$$H_0 = \omega_c a^\dagger a + \omega_m b^\dagger b$$

Optomechanical coupling:

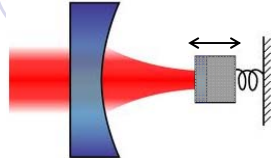
$$g = [\partial \omega_c(x) / \partial x] x_{\text{ZPF}}$$

Coherent laser input


$$H_{\text{drive}} = \Omega^* e^{i\omega_m t} a + \Omega e^{-i\omega_m t} a^\dagger \quad \Omega = \sqrt{\kappa_{\text{ex}} P / (\hbar \omega_m)} e^{i\phi}$$

Linearization $a \rightarrow a_1 + \alpha, b \rightarrow b_1 + \beta$

$$H_L = -\Delta' a_1^\dagger a_1 + \omega_m b_1^\dagger b_1 + (G a_1^\dagger + G^* a_1)(b_1 + b_1^\dagger)$$

$$G = \alpha g \quad \Delta' = \omega - \omega_c + 2|G|^2 / \omega_m$$


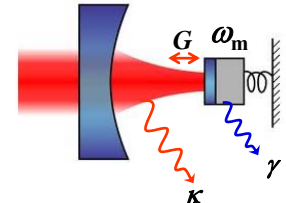
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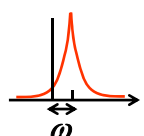
Optomechanical coupling

Important parameters

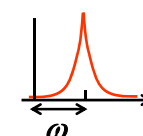
- ω_m : mechanical frequency
- G : coupling strength
- κ : optical decay
- γ : mechanical decay
- n_{th} : thermal phonon number



Unresolved sideband regime: $\omega_m < \kappa$



Resolved sideband regime: $\omega_m > \kappa$



Weak coupling regime: $G < \kappa$

Strong coupling regime: $G > \kappa$

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Cooling and amplification

$$H_L = -\Delta' a_1^\dagger a_1 + \omega_m b_1^\dagger b_1 + (G a_1^\dagger + G^* a_1)(b_1 + b_1^\dagger)$$

Red detuning

Cooling

Blue detuning

Amplification

$H \approx \hbar G(\hat{a}_s^\dagger \hat{b} + \hat{a}_s \hat{b}^\dagger)$

$H \approx \hbar G(\hat{a}_s^\dagger \hat{b}^\dagger + \hat{a}_s \hat{b})$

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Laser cooling of mechanical resonators

Red detuned laser input, imbalanced Stokes & anti-Stokes sidebands.

Anti-Stokes >> Stokes
→ **Cooling but with limit**

Cooling limit (resolved sideband)


Theories:
Wilson-Rae *et al*, *PRL* 99, 093901 (2007)
Marquardt *et al*, *PRL* 99, 093902 (2007)

Experiments:
Nature 444,67 (2006); Nature 444, 71 (2006)
Nat. Phys.4 , 415 (2008); Nat. Phys.5 , 485 (2009)
Nat. Phys.5, 489 (2009); Nat. Phys.5, 509 (2009)
Nature 463,72 (2010); Nature 475, 359 (2011)
Nature 478, 89 (2011); Nature 482 , 63 (2012)

Weak coupling ($G < \kappa$):
Optomechanical coupling is considered as perturbation.

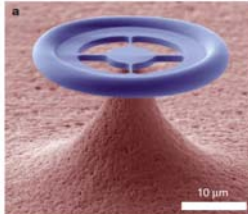
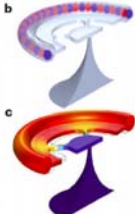
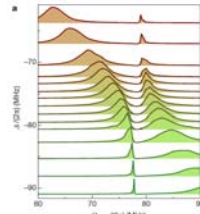
$H_1 = (G a_1^\dagger + G^* a_1)(b_1 + b_1^\dagger)$

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Cooling beyond the weak coupling regimes


From weak to intermediate or even strong coupling regime ($G > \kappa$)

Theory: Dobrindt *et al*, PRL 101, 263602 (2008);
 Experiment: Verhagen *et al*, Nature (2012); Palomaki *et al*, Nature (2013).

- Important for quantum applications
- Requires non-perturbative calculation

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Cooling beyond the weak coupling regimes

Non-perturbative approach:

$$\dot{\rho} = i[\rho, H_L] + \kappa \mathcal{D}[a_1]\rho + \gamma(n_{\text{th}} + 1)\mathcal{D}[b_1]\rho + \gamma n_{\text{th}}\mathcal{D}[b_1^\dagger]\rho$$

$$\dot{\mathbf{V}} = \mathbf{M}\mathbf{V} + \mathbf{N},$$


$$\mathbf{V} = (\bar{N}_a, \bar{N}_b, \langle a_1^\dagger b_1 \rangle, \langle a_1 b_1^\dagger \rangle, \langle a_1 b_1 \rangle, \langle a_1^\dagger b_1^\dagger \rangle, \langle a_1^2 \rangle, \langle a_1^{\dagger 2} \rangle, \langle b_1^2 \rangle, \langle b_1^{\dagger 2} \rangle)^T$$

$$\mathbf{N} = (0, \gamma n_{\text{th}}, 0, 0, -iG, iG^*, 0, 0, 0, 0)^T$$

- Valid in weak, intermediate and strong coupling regimes
- Can study dynamics behavior, in addition to steady state

Y.-C. Liu et al., PRA 89, 053821 (2014)


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Outline

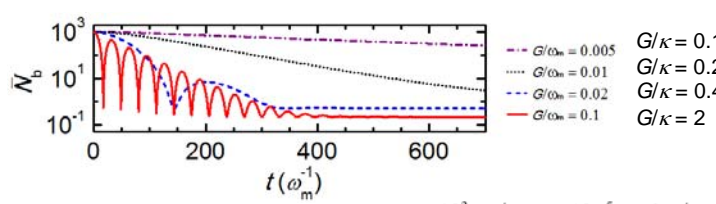
1. Introduction
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Cooling in the strong coupling regime

Time evolution of average phonon number in weak and strong coupling regimes (numerical result)




$n_{\text{th}} = 10^3$, $\gamma/\omega_m = 10^{-5}$, and $\kappa/\omega_m = 0.05$

Characteristics of cooling in strong coupling regime

- **Oscillation and cooling rate saturation**
- **Instantaneous-state cooling limit**

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Cooling in the strong coupling regime

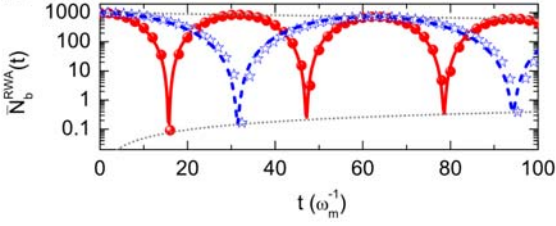
Under rotating-wave approximation (RWA) $|G| \ll \omega_m$

$$\bar{N}_b^{\text{RWA}}(t) \simeq n_{\text{th}} \frac{\gamma + e^{-\frac{\kappa+\gamma}{2}t} [\kappa \cos^2(|G|t) - \gamma \sin^2(|G|t)]}{\kappa + \gamma}$$

$$n_{\text{ins}}^{\text{RWA}} \simeq \frac{\pi\gamma n_{\text{th}}}{4|G|}$$

$t \sim \pi/(2G)$

$n_{\text{th}} = 1000$ $G/\omega_m = 0.1; G/\omega_m = 0.05$




“研究热环境的影响”

$N_b, n_{\text{ins}} \propto n_{\text{th}}$

Data: numerical result;
Curve: analytical result (RWA)

Y.-C. Liu et al., PRA 89, 053821 (2014)

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Cooling in the strong coupling regime

Without RWA, $n_{\text{th}} = 0$ “研究反旋波的影响”


$$\bar{N}_b^{(0)}(t) \simeq \frac{|G|^2 [1 - e^{-\frac{\kappa+\gamma}{2}t} \cos(\omega_+ + \omega_-)t \cos(\omega_+ - \omega_-)t]}{2(\omega_m^2 - 4|G|^2)}$$

$$\omega_{\pm} = \sqrt{\omega_m^2 \pm 2|G|\omega_m}$$

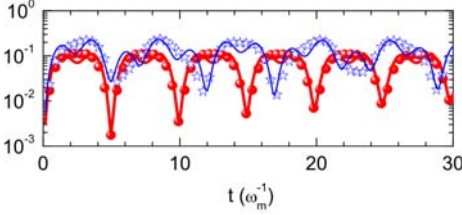
Frequency matching condition:

$(\omega_+ + \omega_-)t = p\pi,$
 $(\omega_+ - \omega_-)t = q\pi,$ both odd or even

p, q

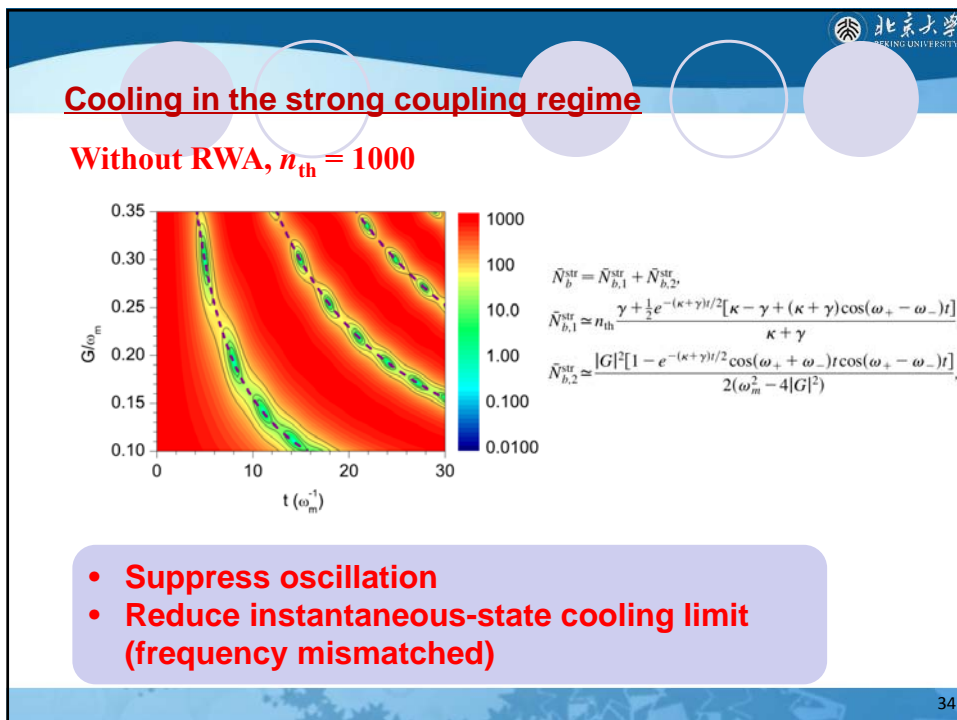
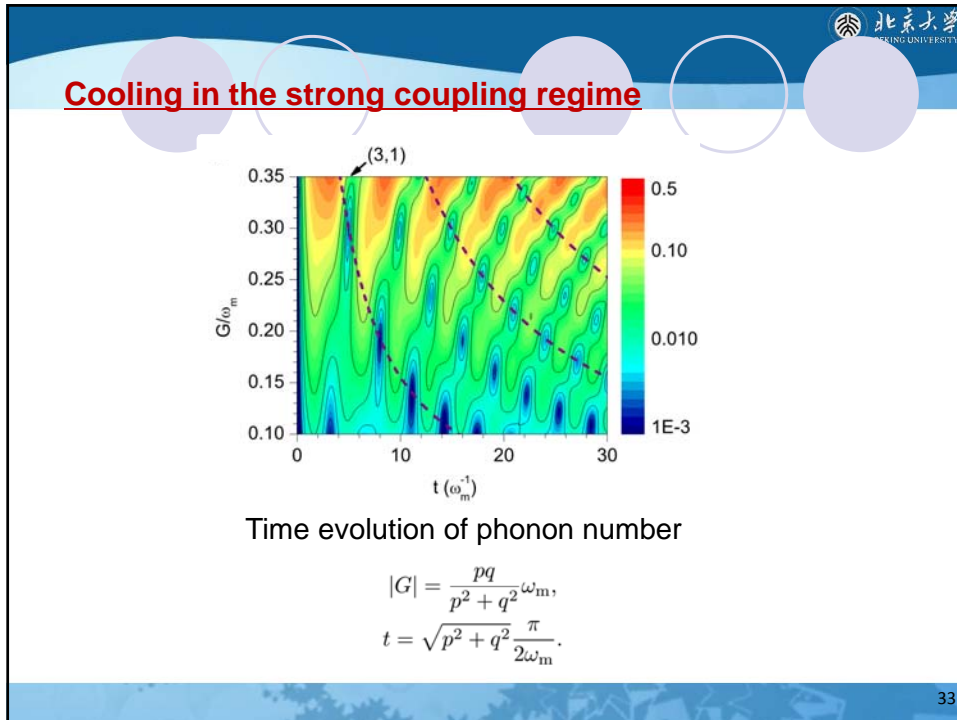


$$n_{\text{ins}}^{(0)} \simeq \frac{\pi\kappa|G|}{8(\omega_m^2 - 4|G|^2)}$$



$G/\omega_m = 0.3$
 $G/\omega_m = 0.35$

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Dynamic dissipative cooling approach

Cooling processes:

- A: Energy swapping
- C: Counter-rotating-wave interaction
- E: Cavity dissipation

Heating processes:

- B: Swap heating
- D: Quantum backaction heating
- F: Thermal heating

Dissipation pulses

n: photon number
m: phonon number

Strong coupling: E is slower than A/B .

➔ Dynamically modulate E process.

Y.-C. Liu *et al.*, *Phys. Rev. Lett.* 110, 153606 (2013)

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Fast cooling to the steady-state limit


A single dissipation pulse

$G/\omega_m=0.2$
 $\kappa/\omega_m=0.05$

Blue: w/o dissipation pulse
Red: with dissipation pulse

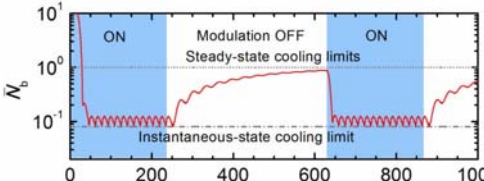
Modulation scheme and mean phonon number.

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Breaking the cooling limit @ frequency mismatched

Periodic dissipation pulses Period: $\pi/(\omega_+ - \omega_-)$



$G/\omega_m = 0.1$
 $\kappa(0)/\omega_m = 0.01$

Instantaneous-state cooling limit:

$$\bar{N}_{ins} \approx \frac{\pi \gamma n_{th}}{4|G|} + \frac{\pi^2 |G|^4}{(\omega_m^2 - |G|^2)(\omega_m^2 - 4|G|^2)}$$


经典 量子

Modulation is switchable.
ON: Instantaneous-state limit
OFF: Steady-state limit

Steady-state cooling limit:

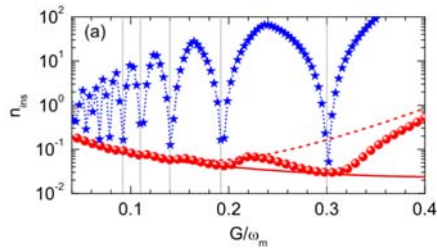
$$\bar{N}_{std}^{str} \approx \gamma n_{th} / (\kappa + \gamma) + |G|^2 / [2(\omega_m^2 - 4|G|^2)]$$

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Breaking the cooling limit @ frequency mismatched

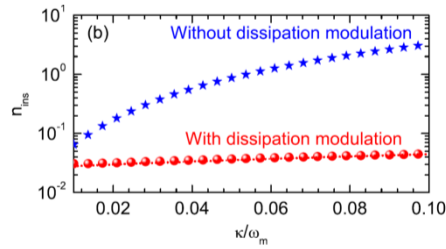
Comparison of cooling limits



声子数 vs. 耦合强度(G)


— Frequency matched
 - - - Frequency mismatched

蓝色: 无周期性调制
 红色: 有周期性调制



声子数 vs. 腔膜耗散(κ)

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Experimental realization

Modulation of cavity decay rate

- **By modulating the free-carrier plasma density.**
 - The absorption is affected by injection of charge carriers into an undoped sample.
- **By using a light absorber/scatterer.**
 - The position change of a light absorber/scatterer leads to the change of the cavity dissipation by absorption of photons and Rayleigh scattering photons out of the cavity.

Q. Xu *et al.*, *Nat. Phys.* **3**, 406 (2007); K. Kondo *et al.*, *Phys. Rev. Lett.* **110**, 053902 (2013); I. Favero *et al.*, *New J. Phys.* **10**, 095006 (2008).


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Cooling in the intermediate coupling regime

Exact results of the steady-state cooling limit:

$$n_s = n_s^{(1)} + n_s^{(0)}$$

Classical cooling limit

Originates from the mechanical dissipation and is proportional to n_{th}

Quantum cooling limit

Originates from the quantum backaction and does not depend on n_{th}

Approximate analytical results under resolved sideband limit:


$$n_s^{(1)} |_{\Delta'=-\omega_m} \simeq \frac{4|G|^2 + \kappa^2}{4|G|^2 \kappa} \gamma n_{th},$$

$$n_s^{(0)} |_{\Delta'=-\omega_m} \simeq \frac{\kappa^2 + 8|G|^2}{16(\omega_m^2 - 4|G|^2)}.$$

ω_m : mechanical frequency
 G : coupling strength
 κ : optical decay
 γ : mechanical decay
 n_{th} : thermal phonon number

at $G \rightarrow 0$, agree to previous result (微扰) Y.-C. Liu et al., unpublished

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Cooling in the intermediate coupling regime

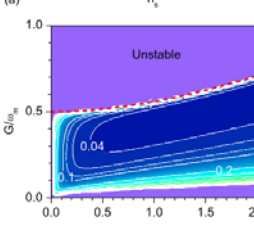
Exact numerical simulation

Classical cooling limit

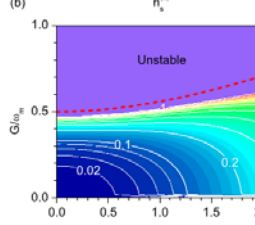
Quantum cooling limit

Total cooling limit

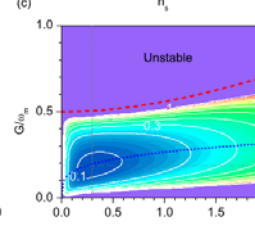
(a) $n_s^{(1)}$



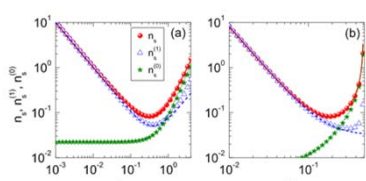
(b) $n_s^{(0)}$




(c) n_s



- Classical cooling limit dominates for small k and G ,
- while quantum limit becomes important as k and G increase.



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Cooling in the intermediate coupling regime

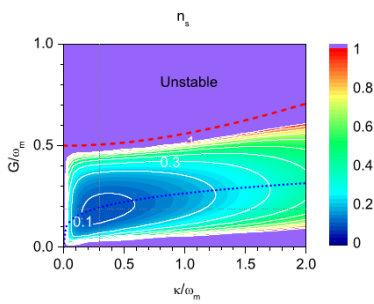
Optimal cooling limit:

$$n_{\text{opt}} \simeq 1.8 \left(\frac{n_{\text{th}}}{Q_m} \right)^{\frac{2}{3}}$$

Optimal parameters:

$$\kappa_{\text{opt}} \simeq 1.5 \omega_m \left(\frac{n_{\text{th}}}{Q_m} \right)^{\frac{1}{3}},$$


$$G_{\text{opt}} \simeq 0.9 \omega_m \left(\frac{n_{\text{th}}}{Q_m} \right)^{\frac{1}{3}},$$

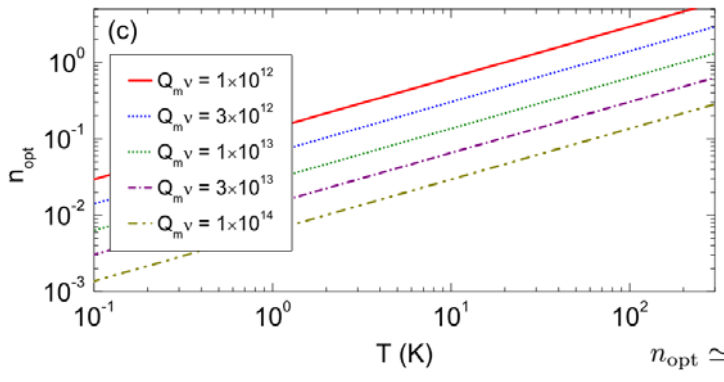


The optimal cooling limit is obtained in the intermediate coupling regime.

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Cooling in the intermediate coupling regime






$n_{\text{opt}} \simeq 1.8 \left(\frac{n_{\text{th}}}{Q_m} \right)^{\frac{2}{3}}$

Ground state: $n_{\text{opt}} < 1 \rightarrow Q_m > 2.4 n_{\text{th}}$

Room temperature: $T = 300\text{K} \rightarrow Q_m \nu > 1.5 \times 10^{13}$

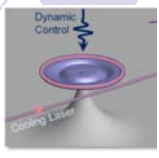
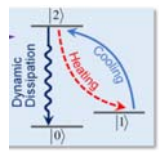
$\nu = 100 \text{ MHz}, Q_m = 1.5 \times 10^5$

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Summary

1. Ultra-high-Q microcavity photonics
2. Cooling in the strong coupling regime
 - Dynamic dissipative cooling
3. Cooling in the intermediate coupling regime
 - Room-temperature ground-state cooling

- Y.-C. Liu, Y.-F. Xiao*, et al., *PRL* 112, 213602 (2014).
- Y.-C. Liu, Y.-F. Xiao* et al., *PRL* 111, 083601 (2013).
- Y.-C. Liu, Y.-F. Xiao* et al., *PRL* 110, 153606 (2013).
- Y.-C. Liu, Y.-F. Xiao* et al., *Science China*, in press
- Y.-C. Liu, Y.-F. Xiao* et al., *PRA* 89, 053821 (2014).
- H.-K. Liu, Y.-F. Xiao* et al., *PRA* 88, 053850 (2013).
- M.-Y. Yan, Y.-F. Xiao* et al., *PRA* 88, 023802 (2013).
- X.-X. Ren, Y.-F. Xiao* et al., *PRA* 87(3), 033807 (2013).
- H.-K. Li, Y.-F. Xiao* et al., *PRA* 85, 053832 (2012)
- Y.-F. Xiao* et al., *PRA* 85, 031805(R) (2012). *Rapid Comm.*
- Y.-C. Liu, Y.-F. Xiao* et al., *PRA* 84, 011805(R) (2011). *Rapid Comm.*
- Y.-C. Xiao* et al., *PRA* 81(5), 053807 (2010).
-

Thank you for your attention!
 For more information:
 Microcavity Photonics and Quantum Optics
<http://www.phy.pku.edu.cn/~yfxiao/>

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