

Few- and Many-body Physics
near Feshbach Resonance:

Taming Multiple Scatterings in A Unitary Gas

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Isszak W. Killam Fund for Advanced Studies

and Institute for Advanced Studies, TsingHua University



Cold-Atom Theory Research in Vancouver

Hyperfine Spin Correlated Physics in Quan. Gases (2003-)

--- Snoek(UvA), Demler (Harvard), Affleck, Semenoff (UBC),
Hui Zhai (IASTU), Junliang Song (Innsbruck), Andrew Ji, W. M.
Liu. (IoP, CAS)...

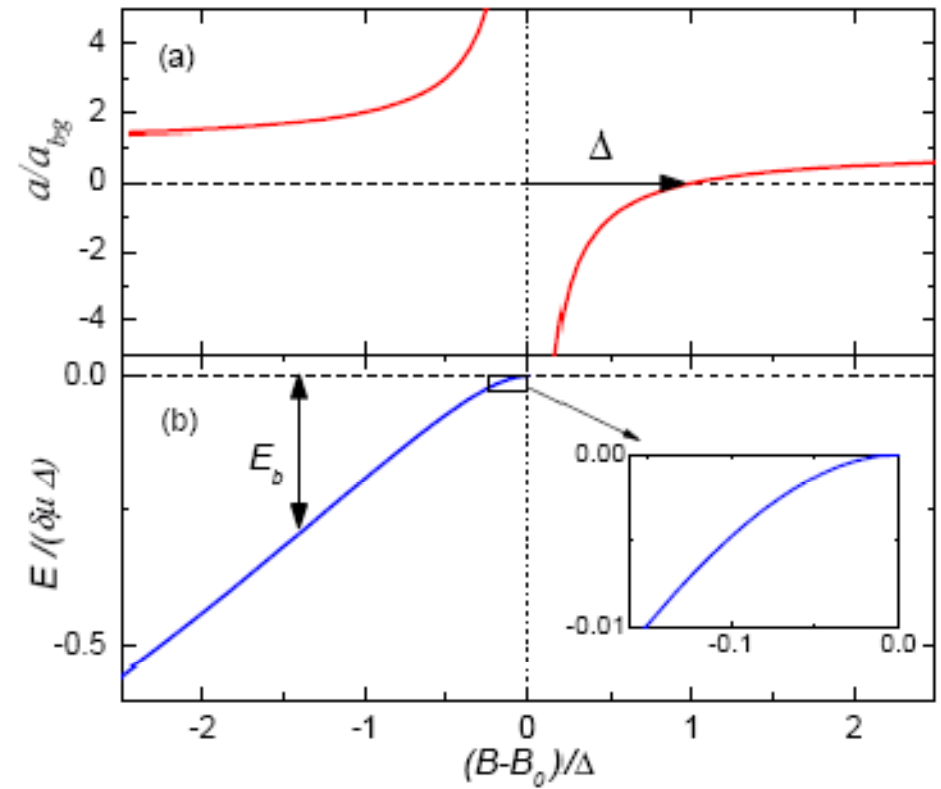
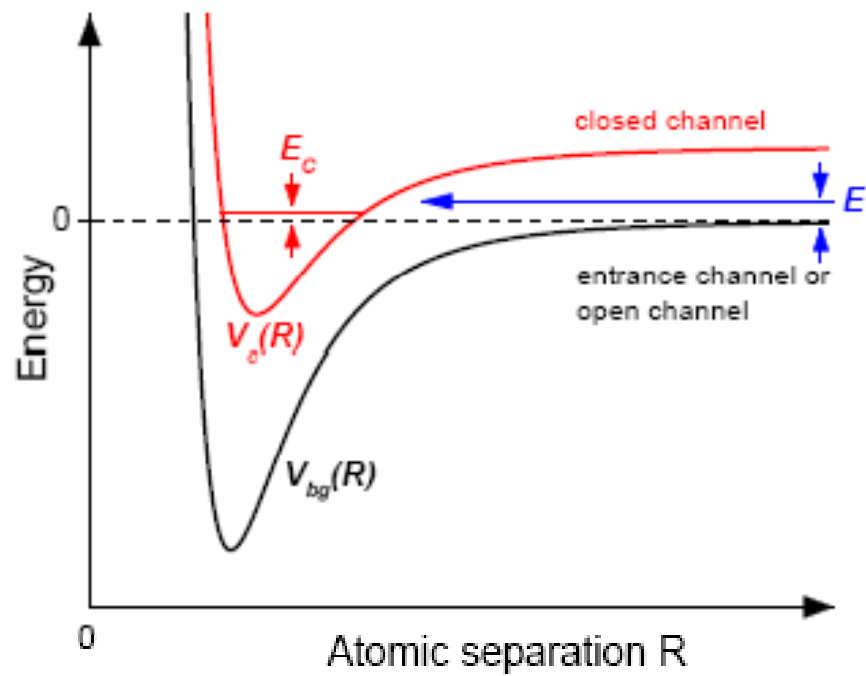
Unitary Gas (2006-)

--- Warner, Zhengcheng Gu (KITP), Alex Huang(UBC), Xiaoling
Cui (IASTU), Junliang Song, David MacNeill(Cornell), Sizhong
Zhang (OSU), Mohammad Mashayekhi, Dmitri Borzov (UBC)

“Spinful” Quantum Gases

- Discrete symmetries; fractionalized vortices and abelian/non-Abelian topological excitations in spinor gases; dynamic generation in rotating traps.
- Field-theoretical models for spin nematic (Z_2 -uniaxial, D_{4h} -biaxial and T -tetrahedral, etc) condensates and Non-MF spin correlated condensates; low-D systems;
- Quantum-fluctuation-induced order from disorder quantum phase transitions, and spin dynamics;

Atomic Feshbach Resonances



detuning

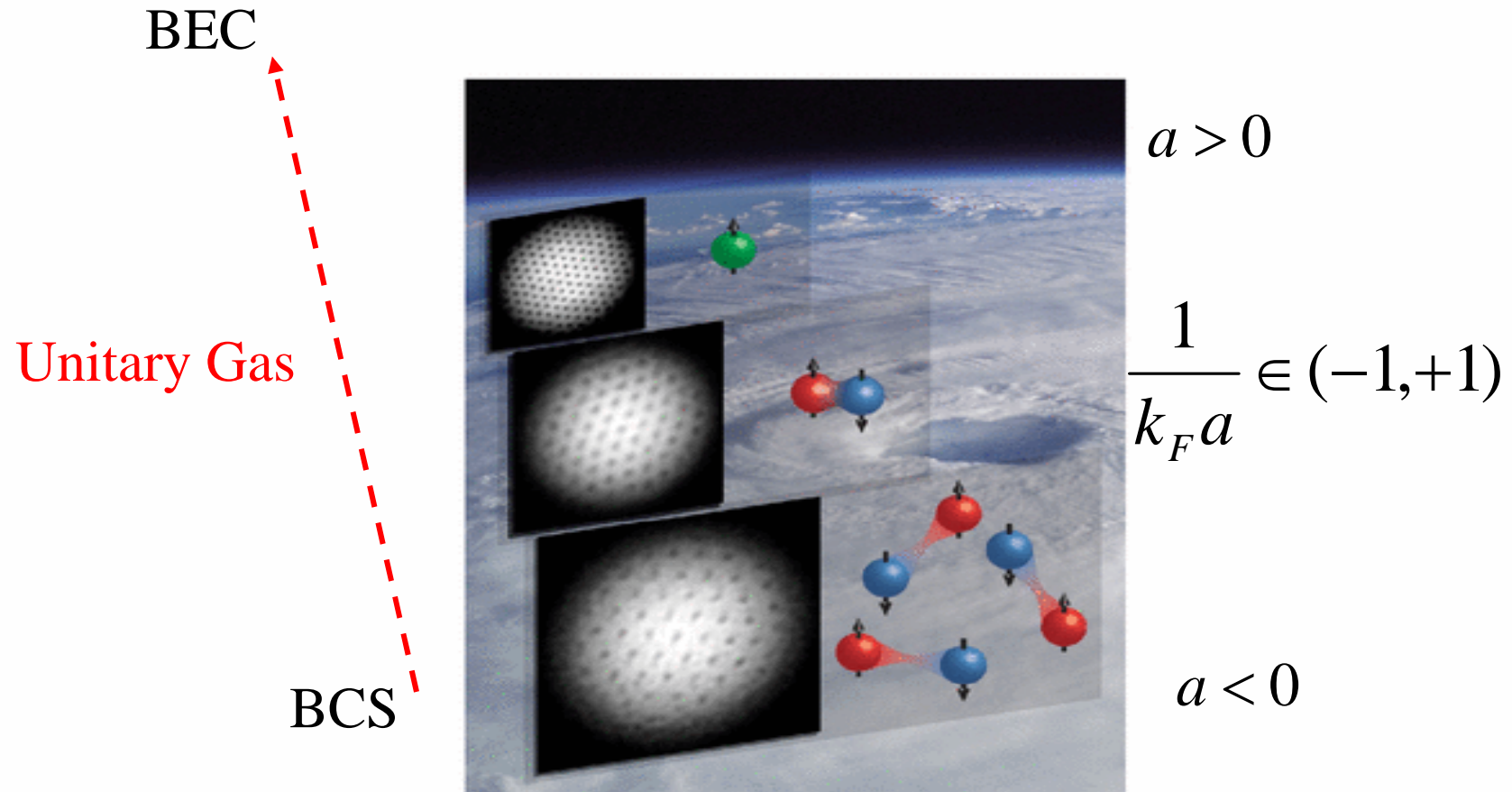


FIGURE 3-5 Vortices in gases. Shown are a vortex pattern in bosonic sodium atoms (green cartoon) in a magnetic trap; vortices in tightly bound lithium molecules (red-blue cartoon); and a vortex lattice in loosely bound fermion pairs on the BCS side of a Feshbach resonance. The background shows a classical vortex (Hurricane Isabel in summer 2003). SOURCE: Wolfgang Ketterle, Massachusetts Institute of Technology.

Outline

Lecture I: Phase shifts in Feshbach resonances: Fundamentals

Phase shifts, scattering lengths and bound states;
Feshbach resonance versus potential scattering;

Lecture II: Resonance Scatterings as a critical phenomenon

RGE approach to resonance: two-body and three-body running
coupling constants; limit cycle behavior;

RGE approach to confinement induced resonance and optical
lattices.

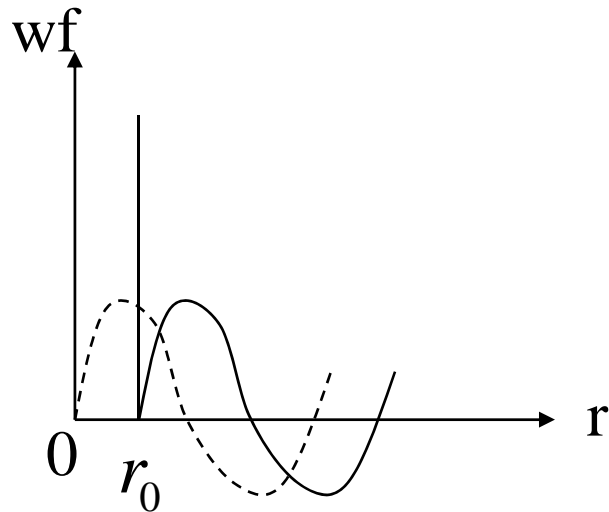
Lecture III: Few-body phenomena in a quantum gas

Many-body effects on few-body scattering phase shifts;
on few-body structures. Dressed few-body clusters.

Lecture IV: Taming Divergence in the Dilute Gas Theory

Nearly fermionized bosons versus Efimov physics;
A diagrammatic approach recently developed in Vancouver
to regularize the theory beyond LYH limit.

S-wave Phase shifts and scattering amplitude



$$\psi_0 = \frac{\sin kr}{r},$$

$$\psi_{hs} = \frac{\sin(kr + \delta)}{r}, \delta = -kr_0$$

$$\psi_{out} \propto e^{ikz} + \frac{f(k)}{r} e^{ikr} = \left(\frac{e^{ikr} - e^{-ikr}}{2kr} \right) + \frac{f(k)}{r} e^{ikr}.$$

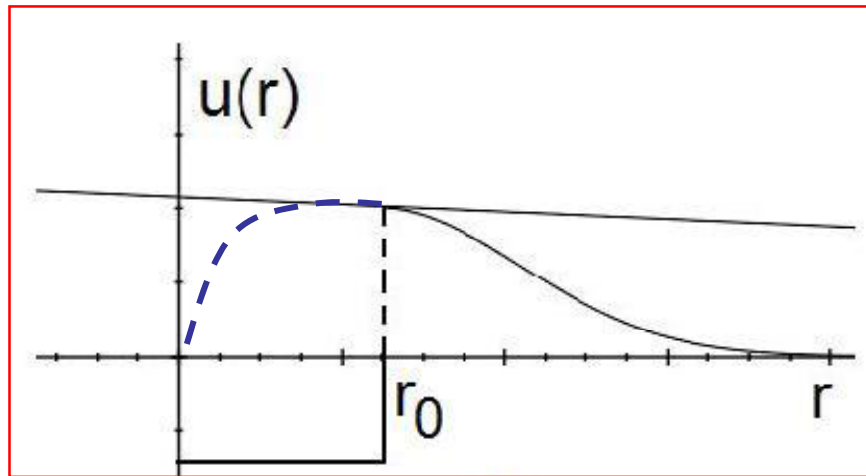
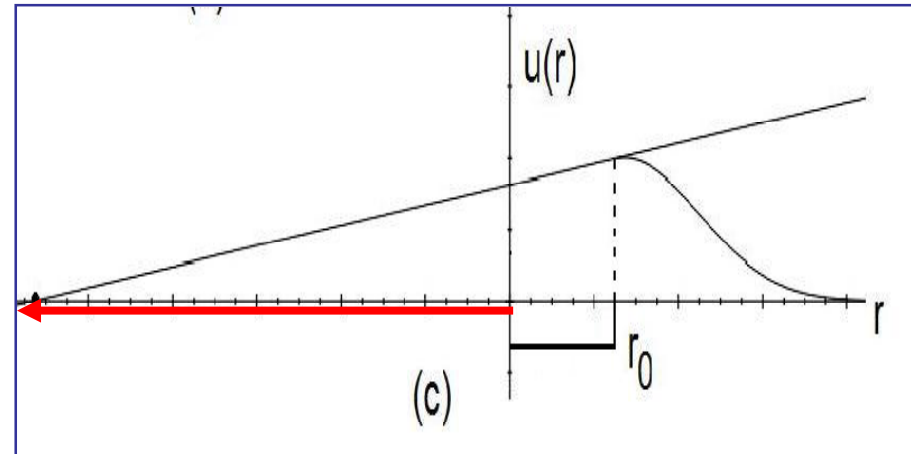
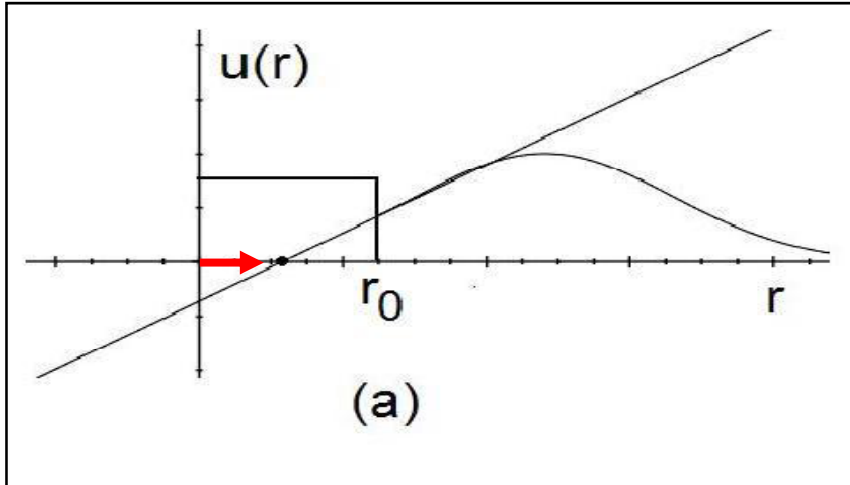
$$f(k) = \frac{e^{i2\delta} - 1}{2ik} = \sin \delta \frac{e^{i\delta}}{k}$$

Phase shifts and scattering lengths: derivatives of the wf outside a potential

$$a = -\lim_{k \rightarrow 0} \frac{\tan \delta(k)}{k}.$$

$$\psi = \frac{\sin(kr + \delta(k))}{r}, \rightarrow \frac{(r\psi)'}{r\psi} = k \lim_{r=0} \frac{\cos(kr + \delta)}{\sin(kr + \delta)} = -\frac{1}{a}$$

Phase shifts, Scattering lengths and bound states



Scatter length diverges when the phase shift becomes $\pi/2$ when a bound state appears.

Scattering length versus amplitude

$$f(k) = \frac{e^{i2\delta} - 1}{2ik} = \sin \delta \frac{e^{i\delta}}{k}$$

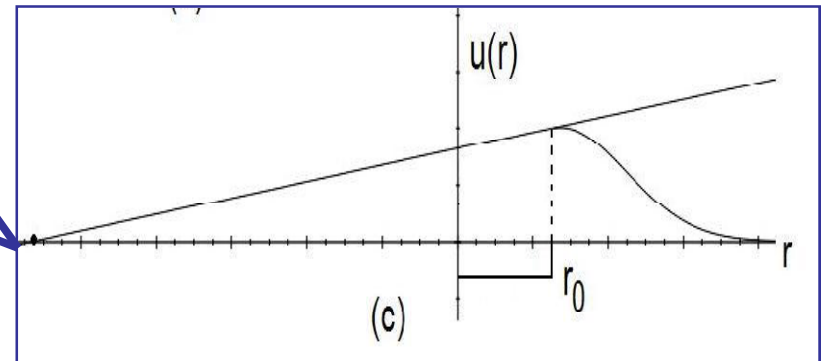
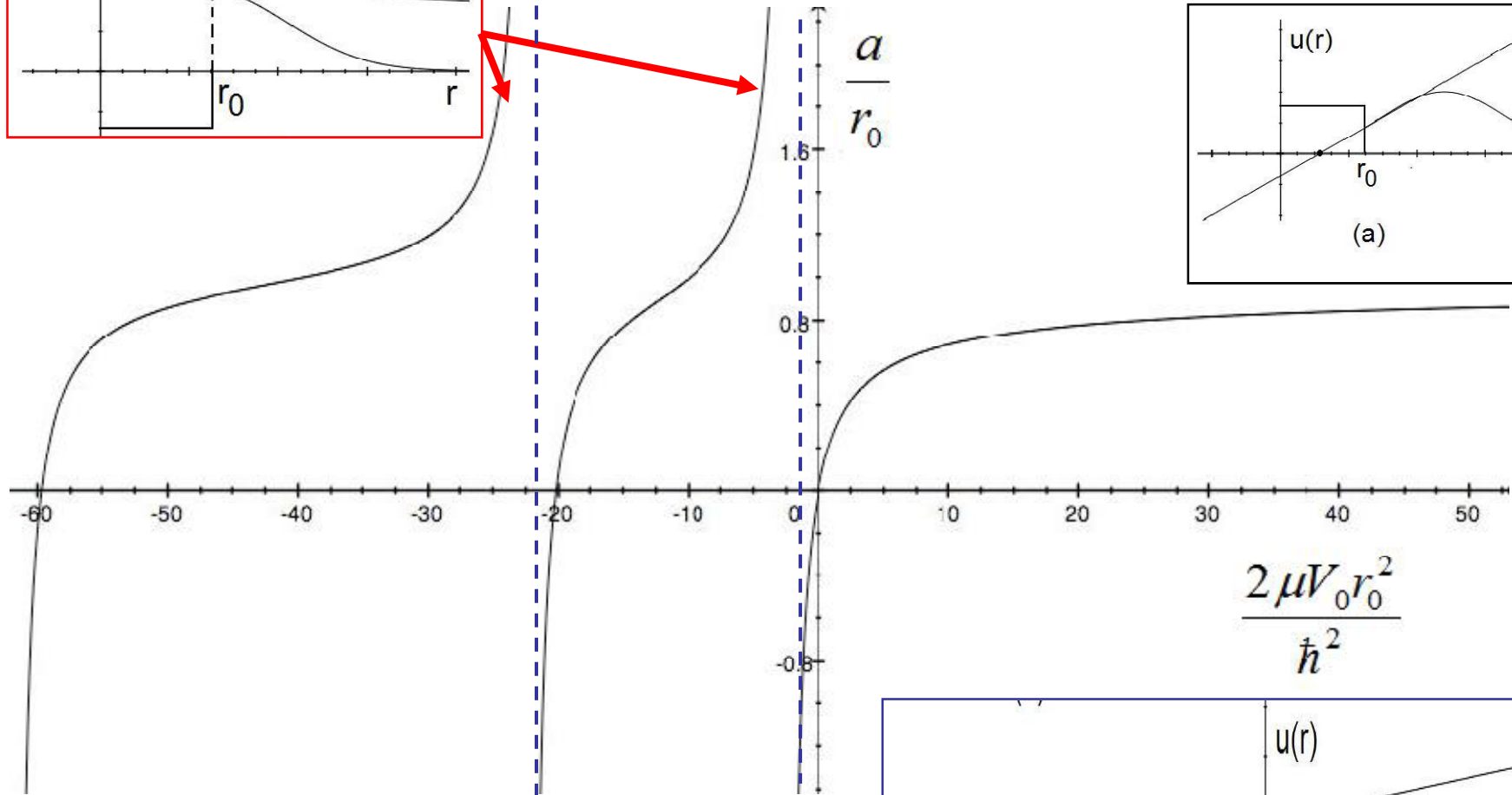
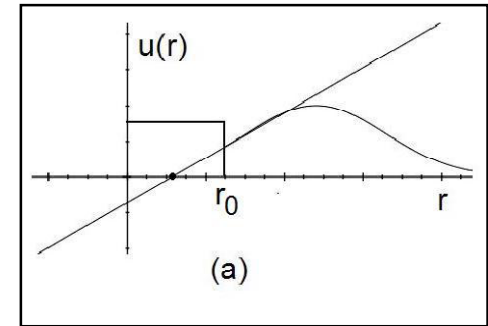
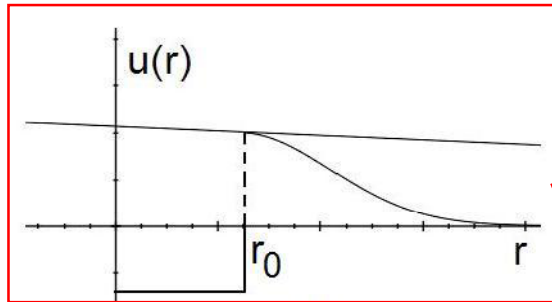
$$f(k) = |f| e^{i\delta(k)} \Leftrightarrow \frac{1}{\frac{1}{a} + ik - C \frac{r_0}{2} k^2}.$$

Scattering cross-section area

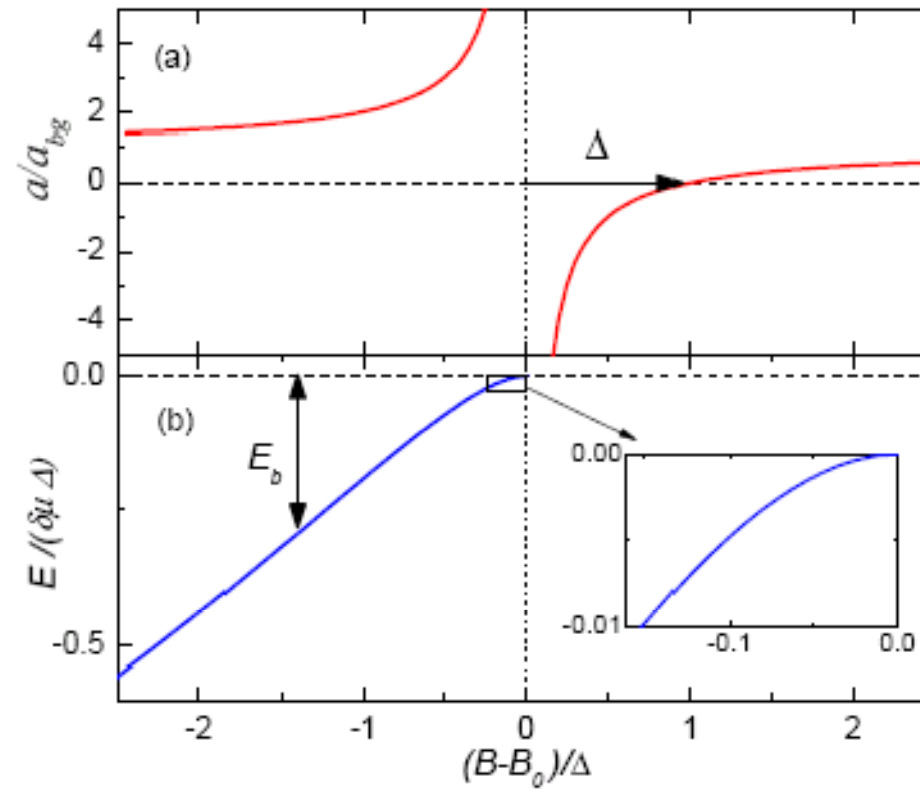
$$\tan \delta(k \rightarrow 0) \approx -ka, \quad \text{or } a = -\lim_{k \rightarrow 0} \frac{\tan \delta(k)}{k}.$$

$$\sigma(r_0 = 0) = 4\pi \frac{\sin^2 \delta(k)}{k^2} = 4\pi \frac{a^2}{1 + k^2 a^2}.$$

Universal Physics for contact interactions



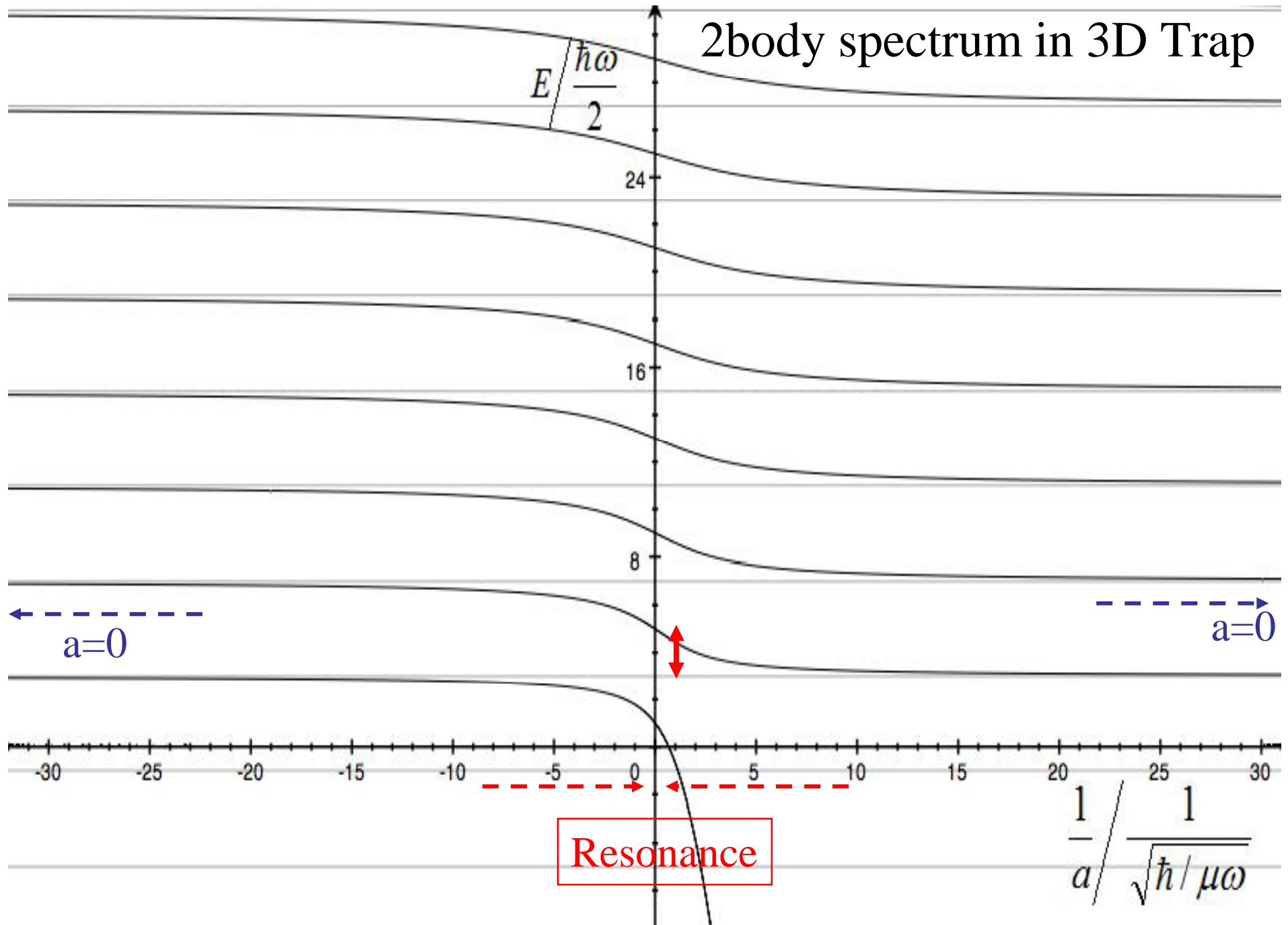
Binding energy versus scattering length



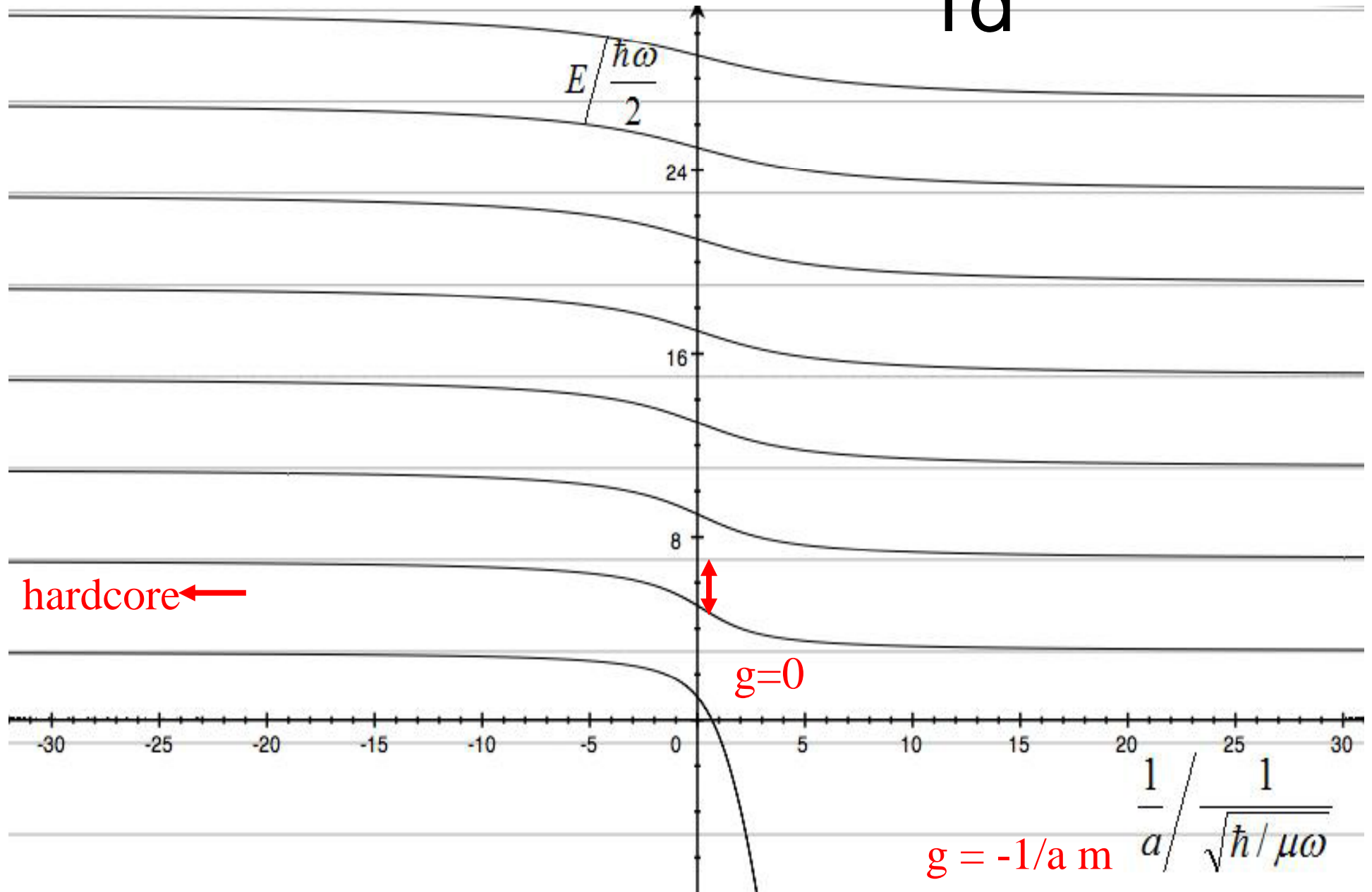
$$E_b = -\frac{1}{2\mu a^2} = -\frac{1}{ma^2}, \quad \mu = \frac{m}{2}.$$

Scaling dimension of interaction energy
versus phase shift/scattering length
(Alex Huang thesis at UBC, 2008)

2body spectrum in 3D Trap



1d



Life time of Feshbach molecules

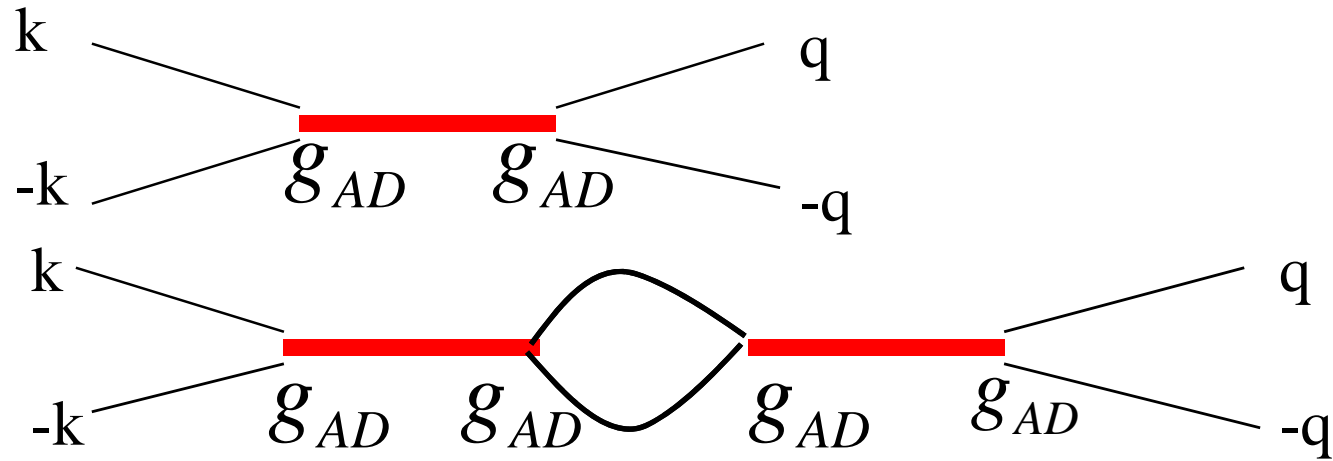


$$G^0(\omega) = \frac{1}{\omega - \delta_T + i\varepsilon} \rightarrow$$

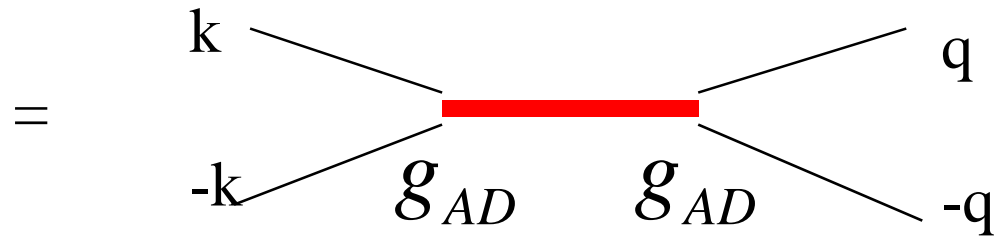
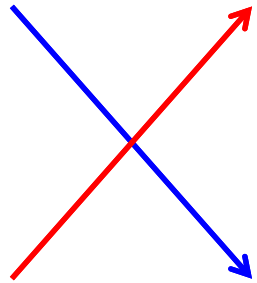
$$G^D(\omega) = \frac{1}{\omega - \left(\delta_T - \frac{mg_{AD}^2 \Lambda^*}{2\pi^2} \right) + im^{3/2} g_{AD}^2 \frac{\omega^{1/2}}{2\pi}}$$

Feshbach Resonance—a diagrammatic S-matrix approach (Dyson series of S-matrix)

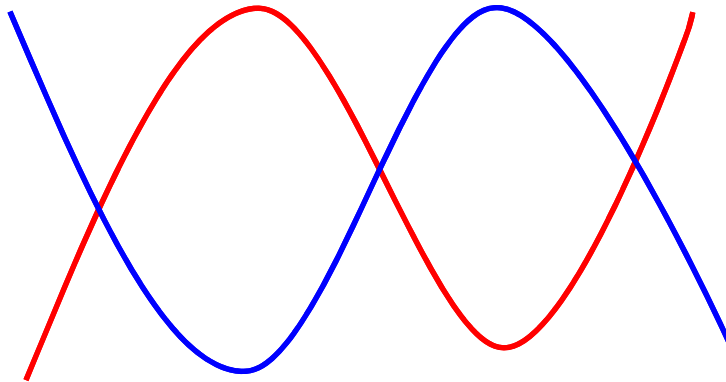
— atom
— dimer



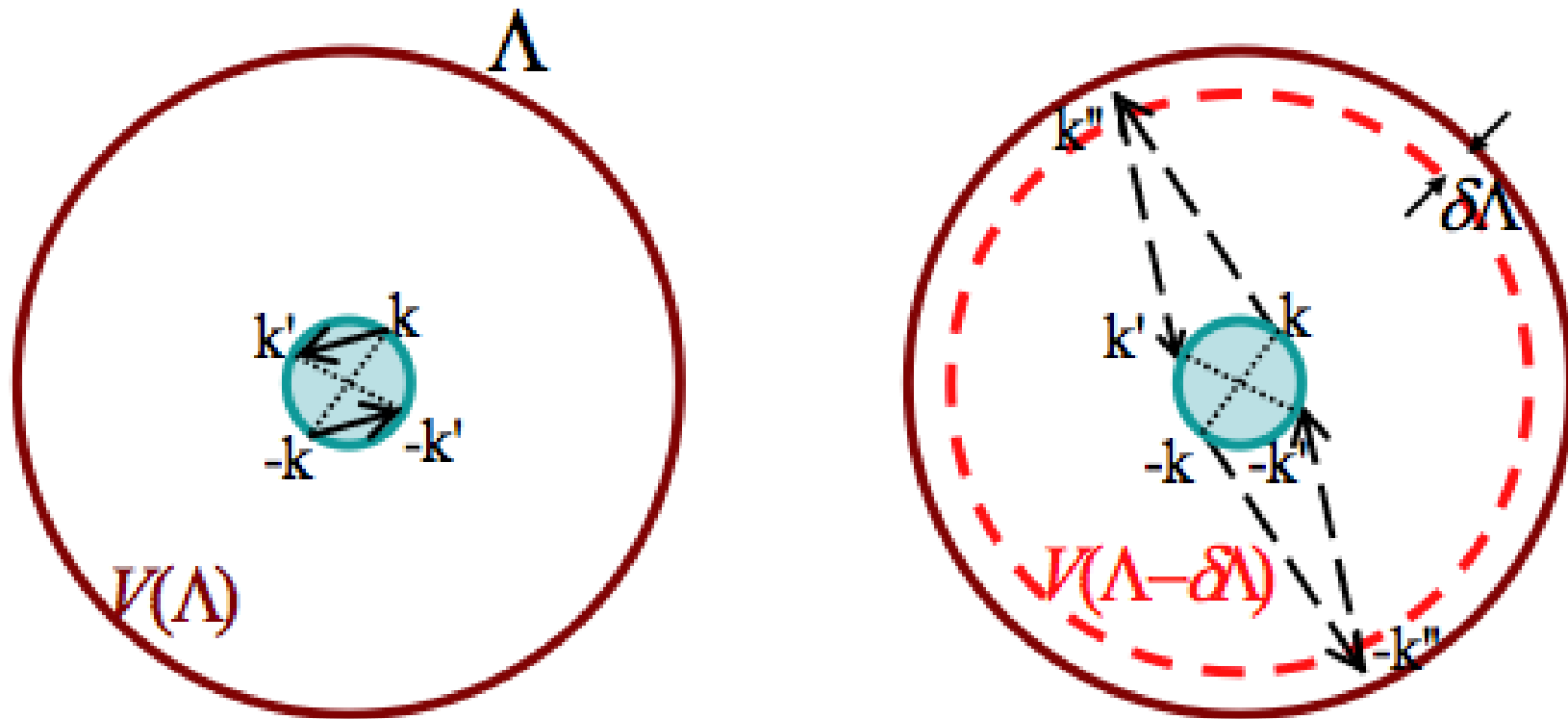
Effective interaction



$$g_{AA}(\omega) = \frac{g_{AD}^2}{\omega - \delta_T + i\delta}$$



Summing up the virtual processes



Properties of Feshbach Resonances

$$f = \frac{1}{\frac{1}{a} + ik - \frac{1}{2} r_e k^2}, \quad a \propto -\frac{m g_{AD}^2}{\delta_D}, \quad r_e \propto -\frac{1}{m^2 g_{AD}^2}.$$

$$[g_{AD}] = [E][V]^{1/2}$$

The scattering amplitude depends on the atom-dimer coupling and detuning which physically is the tunable magnetic field.

When the coupling is large, the scattering amplitude is universal.

Estimate of the effective range

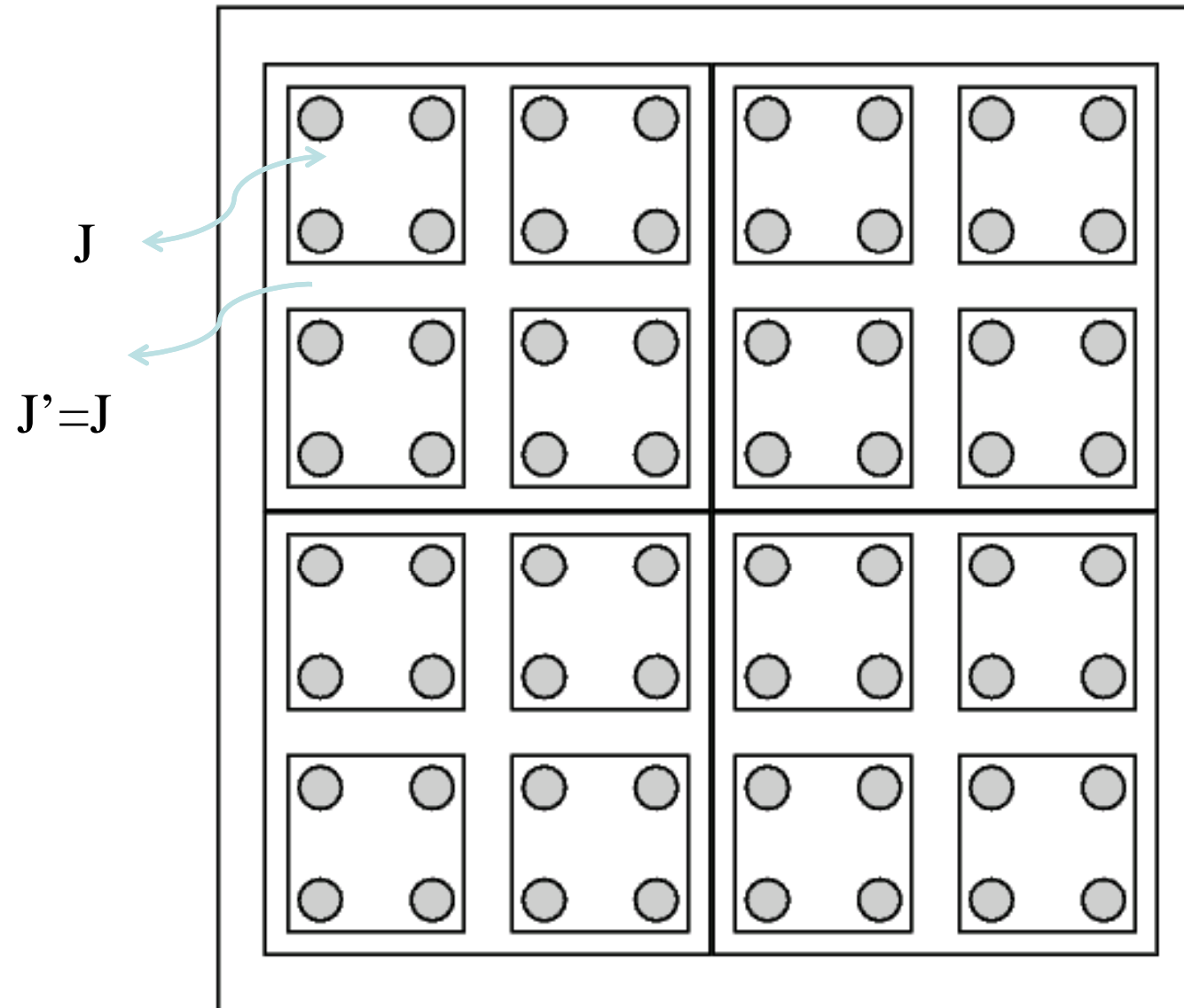
$$a = \left(1 - \frac{\Delta B}{B - B_C}\right) a_{bg} \propto \frac{m g_{AD}^2}{\delta_D}$$

$$\Rightarrow |r_e| \propto \frac{1}{a_{bg} \Delta \mu_B \Delta B m}$$

Lecture II

Resonance scatterings as a critical phenomenon
Infrared fixed points and limit Cycle

Kadanoff scaling at critical points

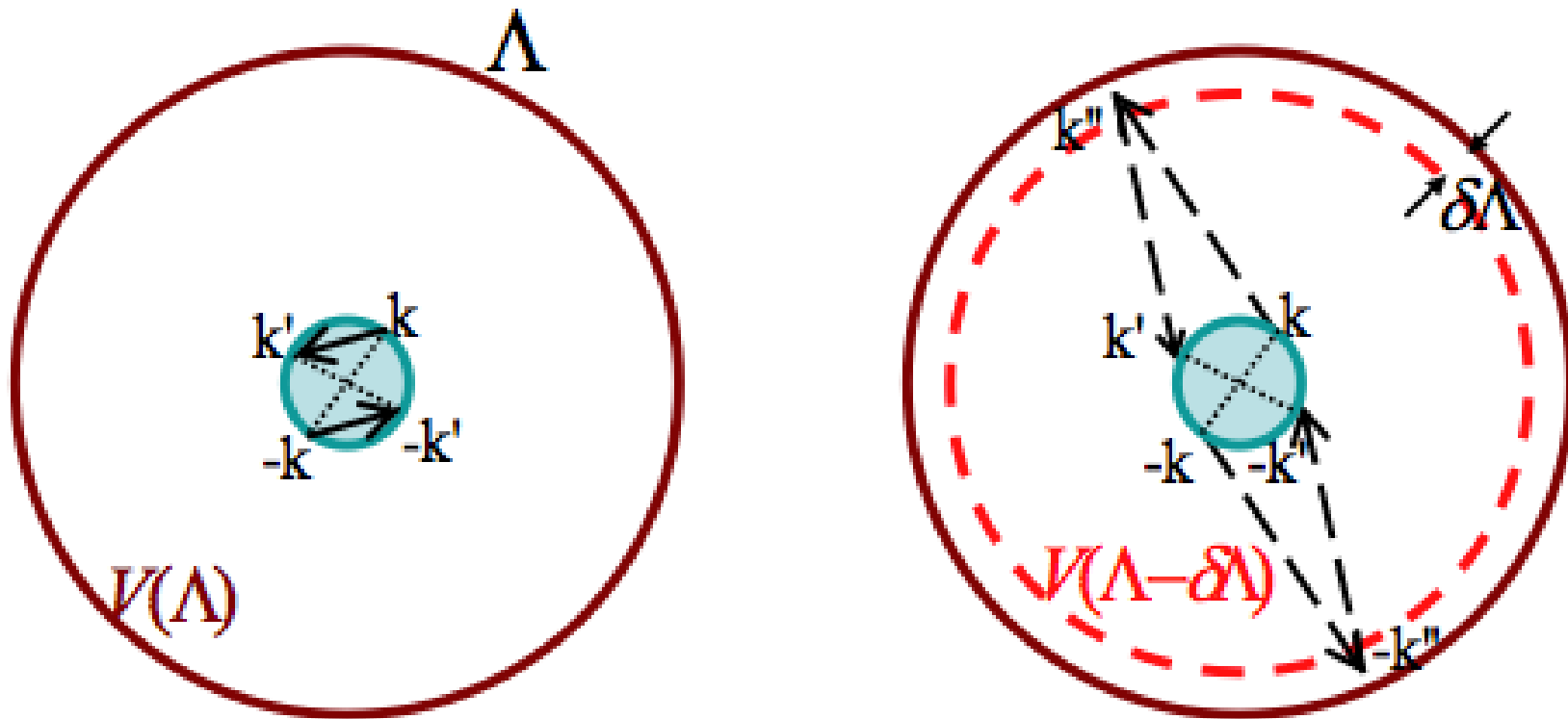


Resonance scattering state

$$\psi(r) = \frac{\sin(kr + \delta)}{kr} \propto 1 - \frac{a}{r};$$

$$r \ll a, \quad \psi(\lambda r) \rightarrow \lambda^{-1} \psi(r)$$

Summing up k-space multiple scatterings



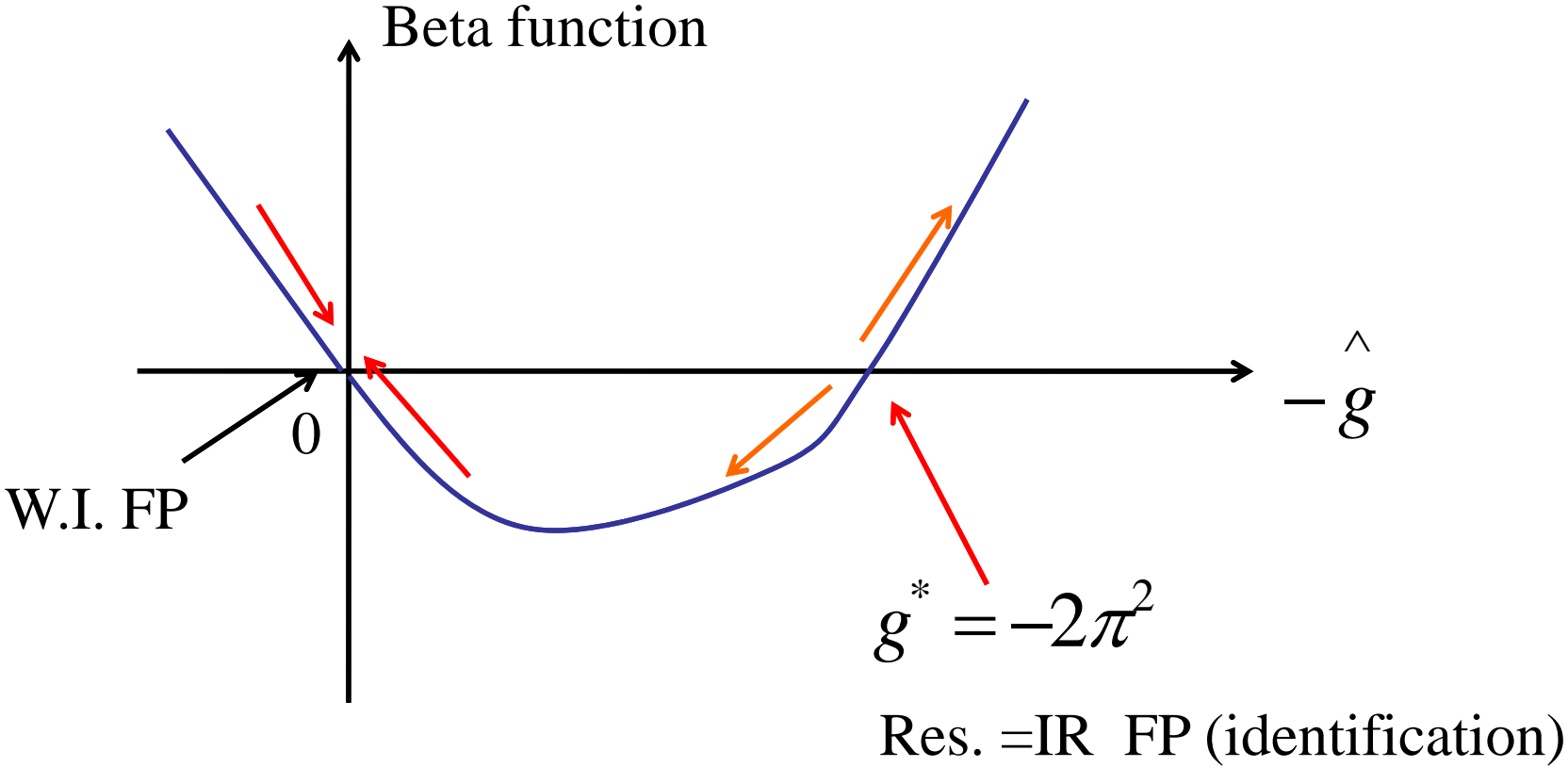
RG equation

(Kaplan et al 98; Nishida and Son et al, 2006;
Nikolic and Sachdev et al, 2007)

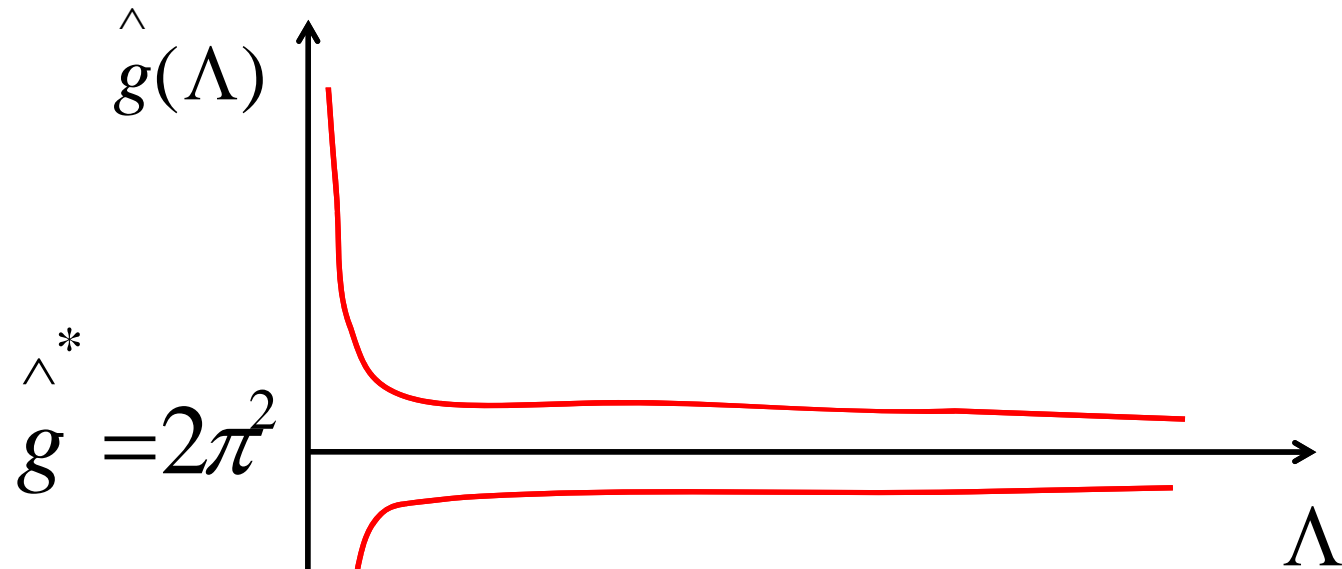
$$\hat{g}(\Lambda) = g_2 * \Lambda,$$

$$\frac{d\hat{g}}{d\log\Lambda} = \beta(\hat{g}), \quad \beta(\hat{g}) = \hat{g} + \frac{1}{2\pi^2} \hat{g}^2$$

Resonance as a unstable FP (3D)



RG flow of the coupling const. and critical exp.



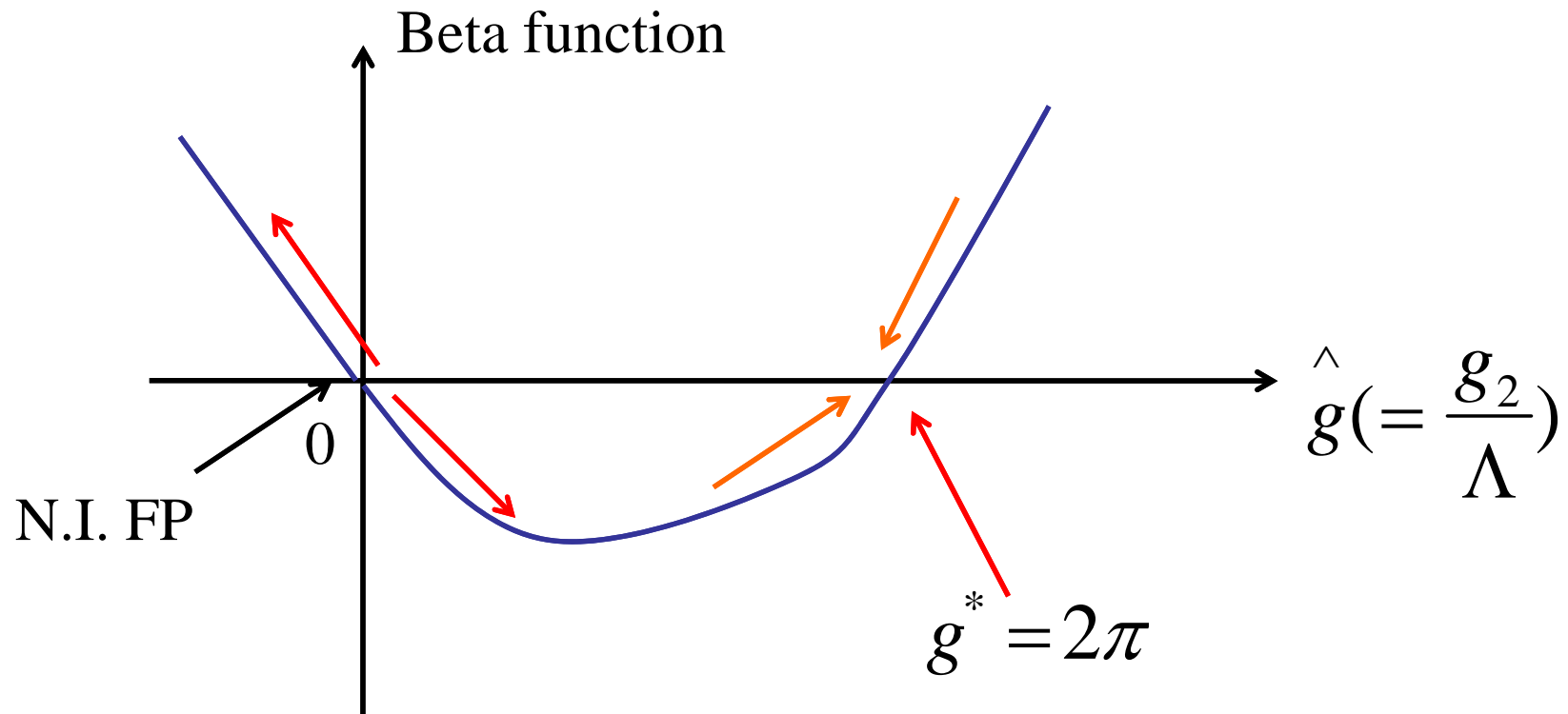
$$a \propto \frac{r_e}{g - g^*}$$

1D

$$\hat{g}(\Lambda) = \frac{g_2}{\Lambda}$$

$$\frac{d\hat{g}}{d\log\Lambda} = \beta(\hat{g}), \quad \beta(\hat{g}) = -\hat{g} + \frac{1}{2\pi}\hat{g}^2$$

Resonance as a unstable FP (1D)



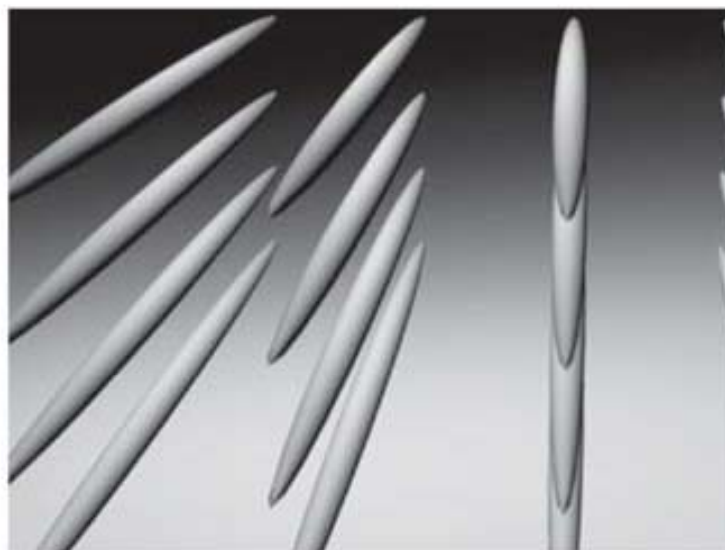
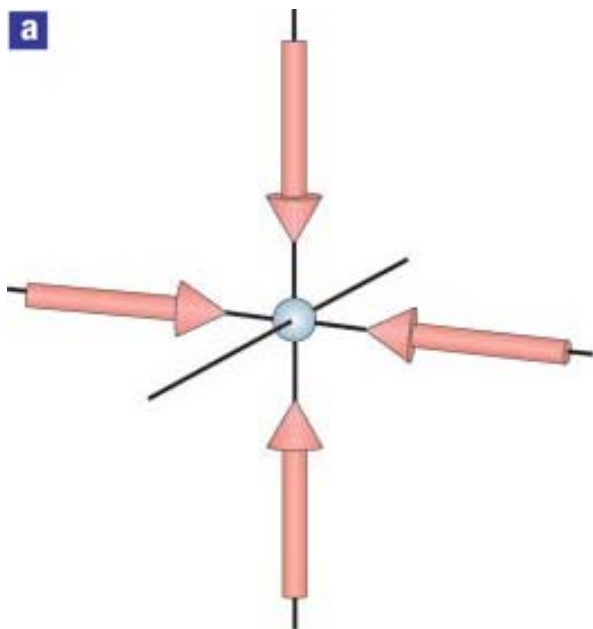
IR FP= Tonk-Girardeau Gas
(Fermionized Bose gas, identification)

Two applications of RG

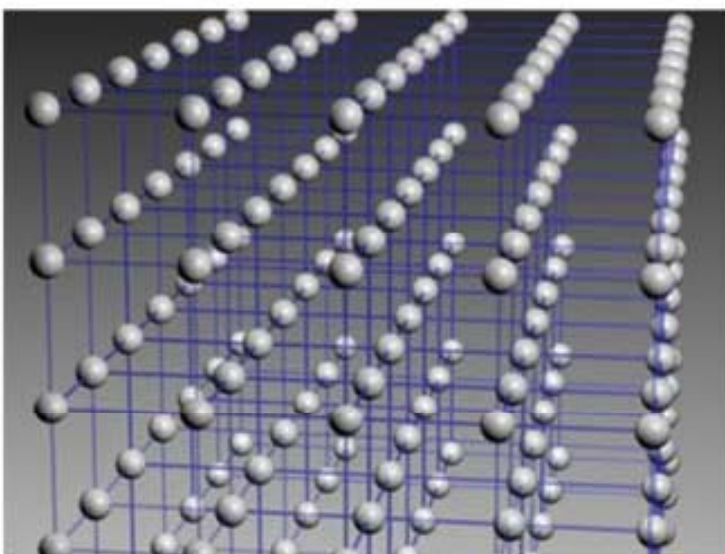
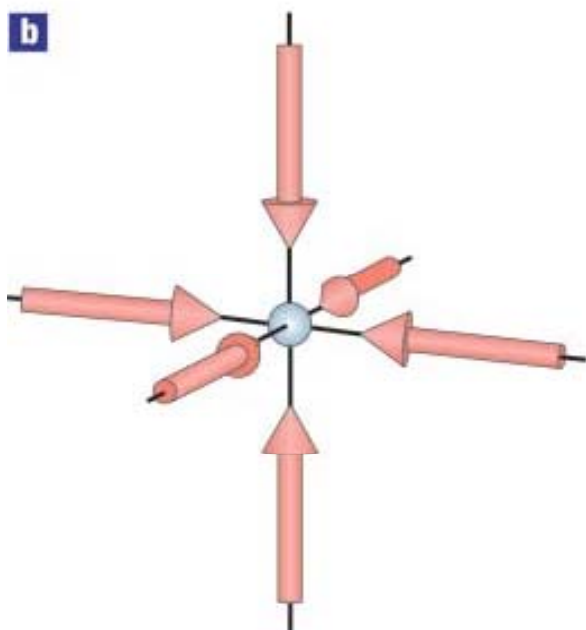
1) Confined geometry:

2) Implications on upper branch physics:

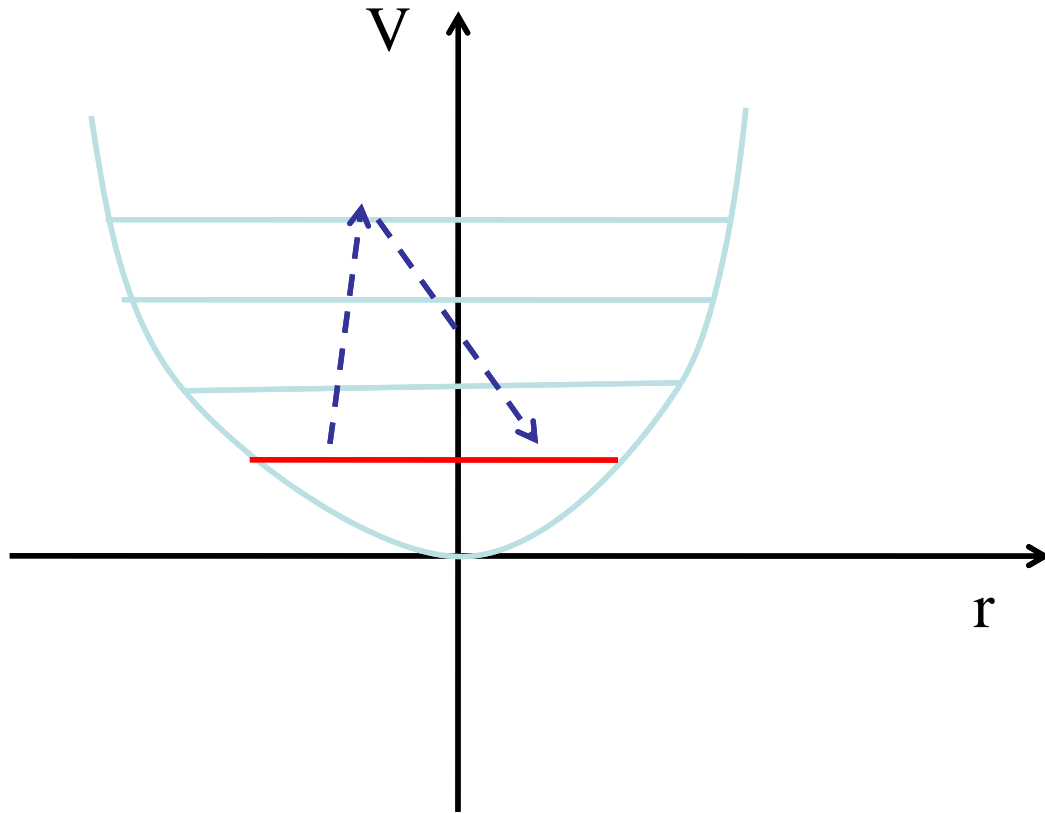
a



b



RG in a confined space---a tube

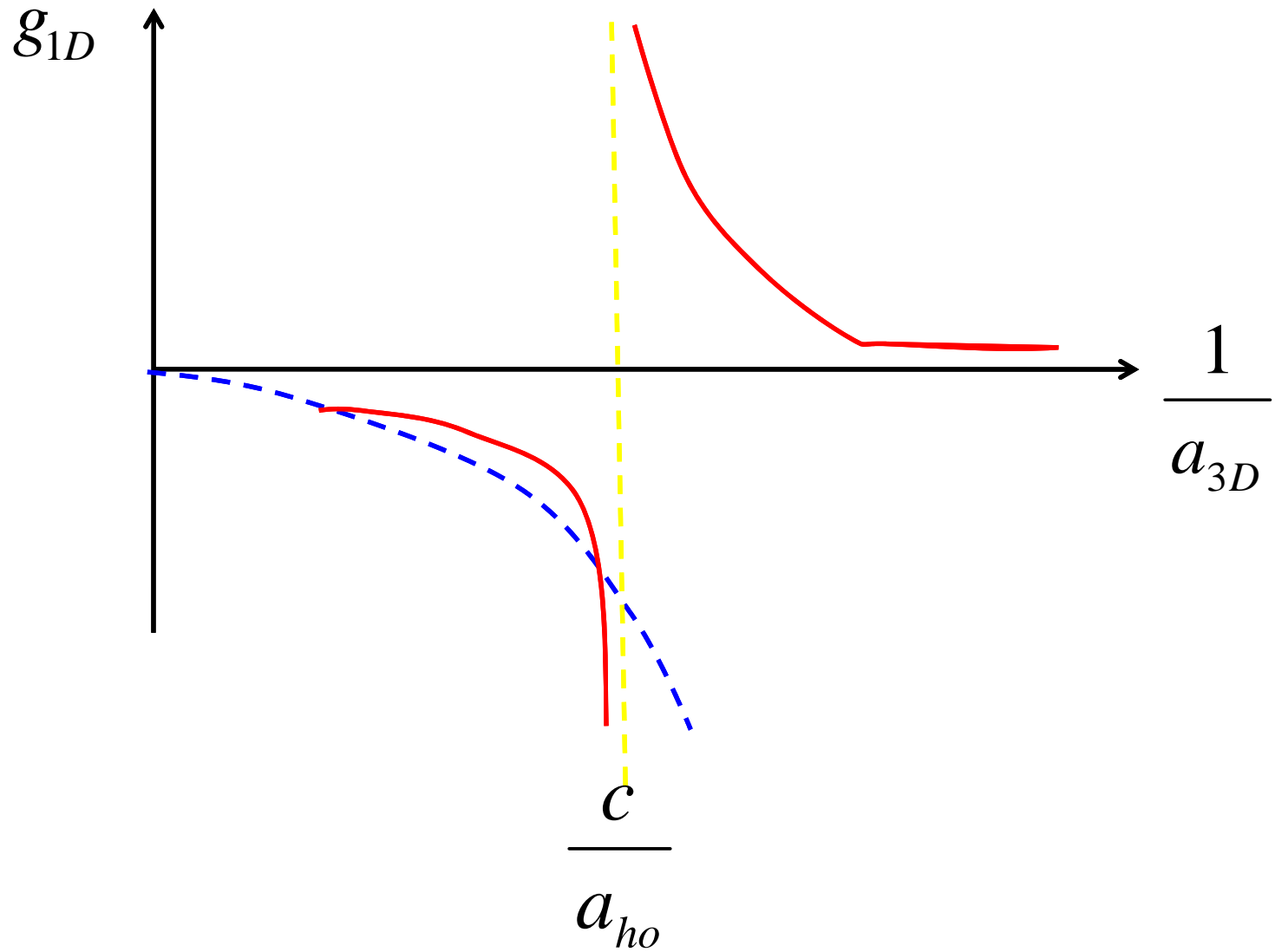


$$g_{1D} \propto \frac{a_{3D}}{a_{ho}^2} \left[\frac{1}{1 - C \frac{a_{3D}}{a_{ho}^2}} \right]$$

$$\varepsilon(n_x, n_y, k_z) = (n_x + n_y + 1)\hbar\omega + \frac{k_z^2}{2m}$$

Olshanii, 98;
Cui, Wang, FZ, 2009.

Confinement induced “hardcore limit”



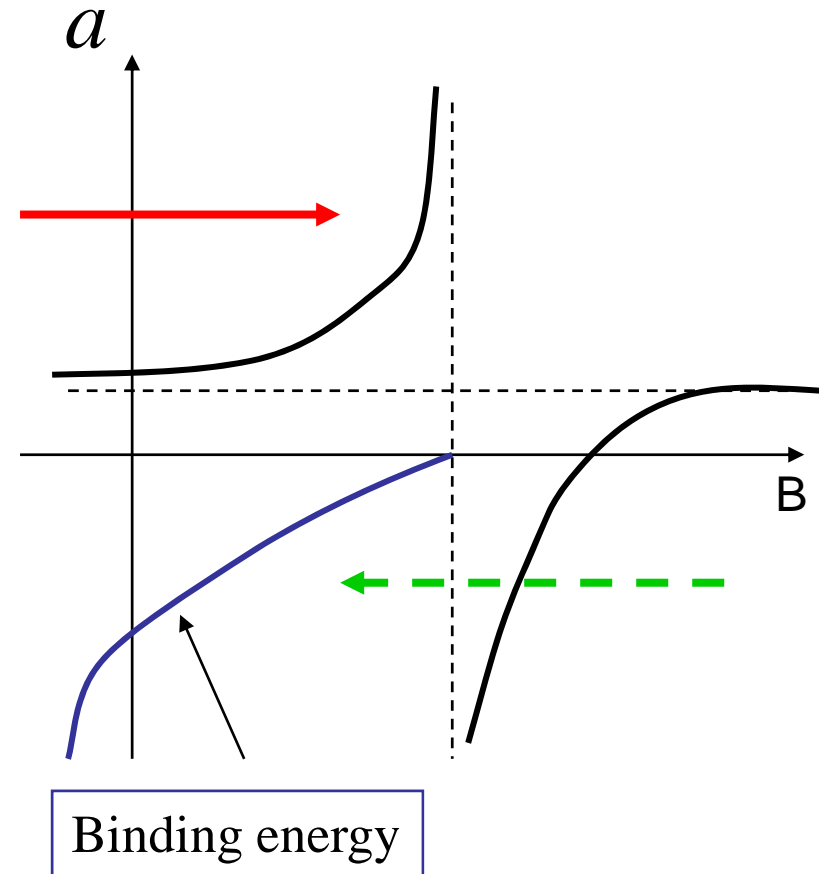
Two applications of RG

1) Confined geometry:

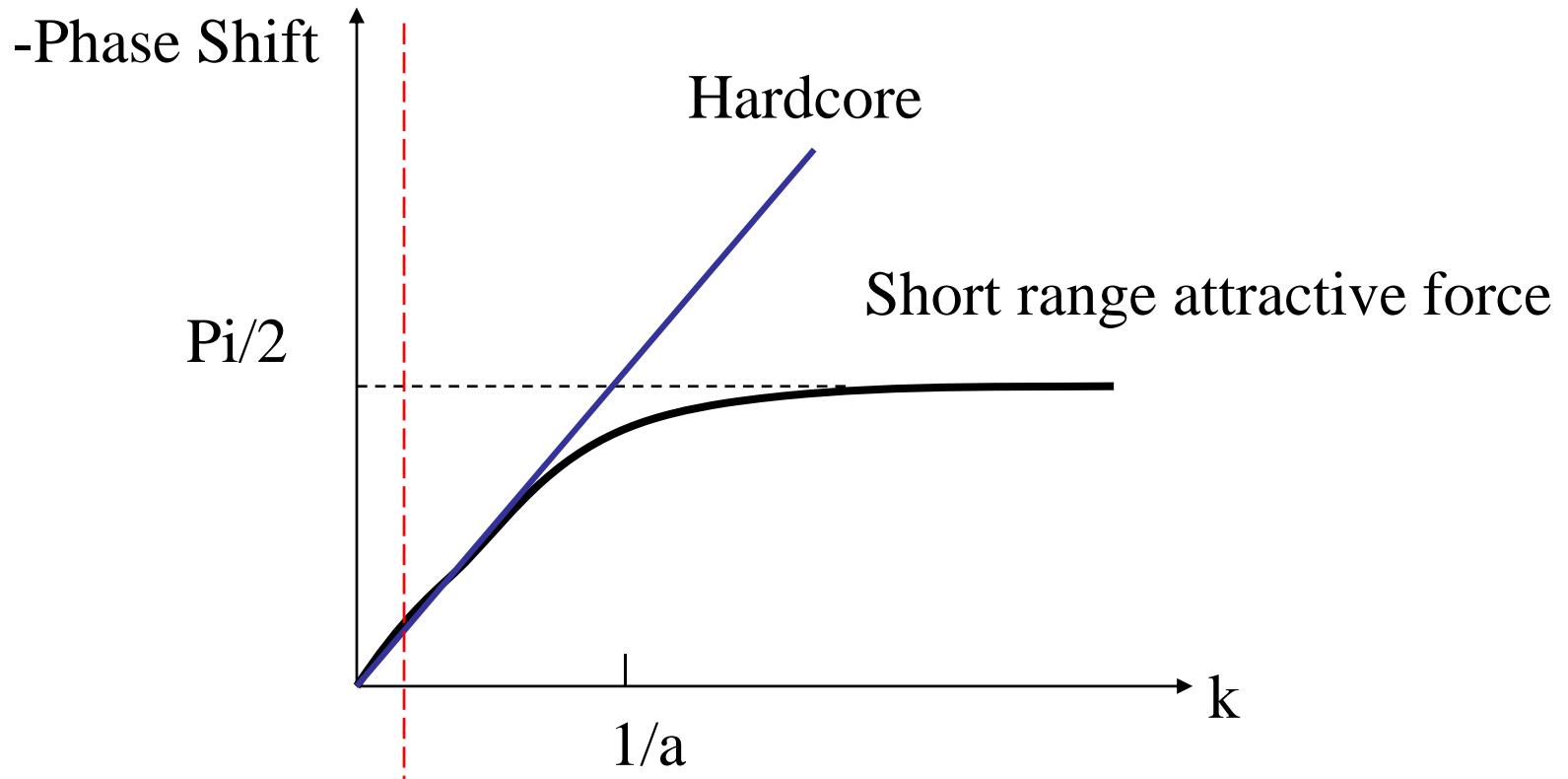
2) Implications on upper branch physics:

Quantum Gas near Feshbach Resonance (Upper Branch)

Is positive “ a ” equivalent to repulsive interactions?

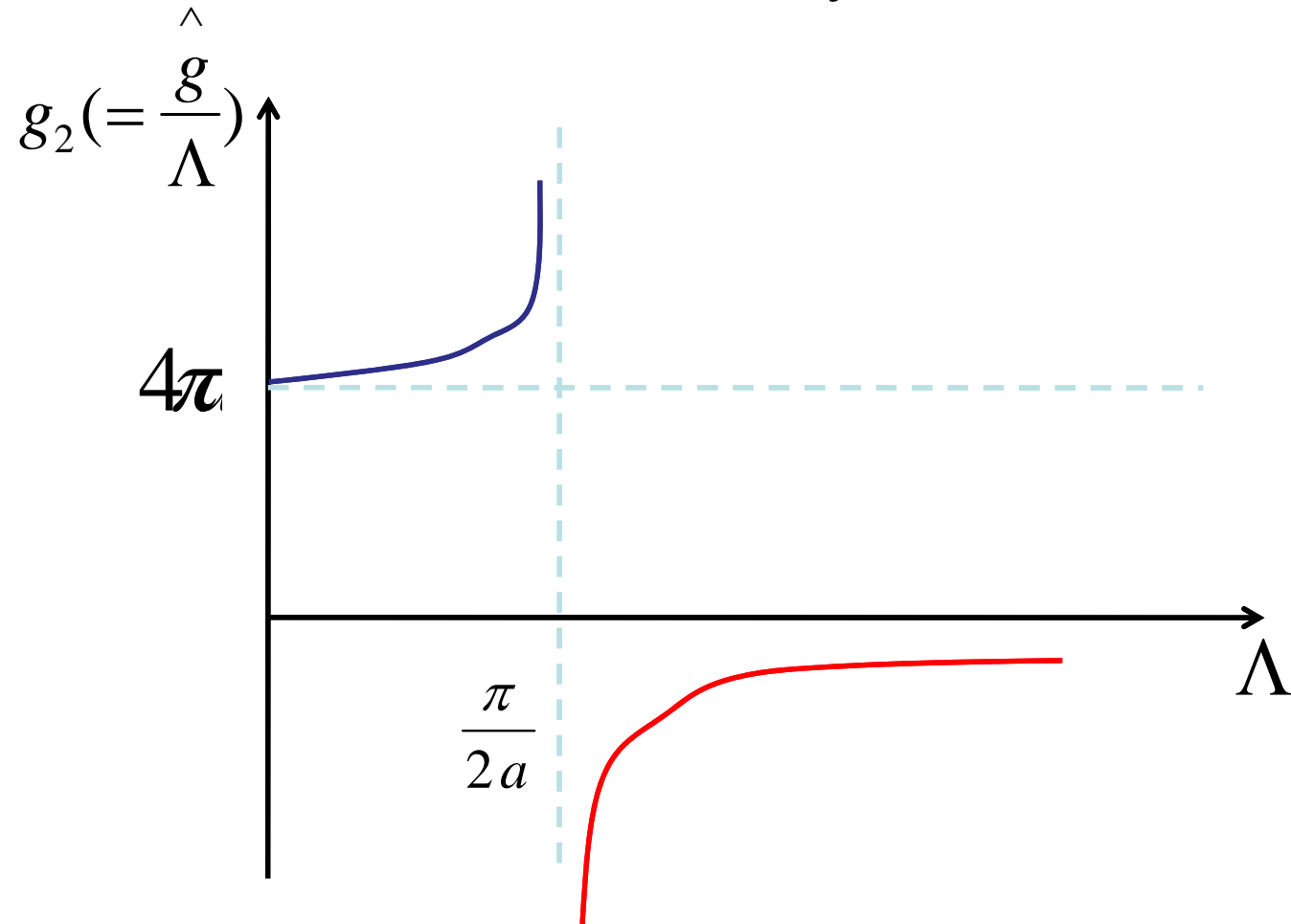


Positive scattering lengths versus hardcore interactions



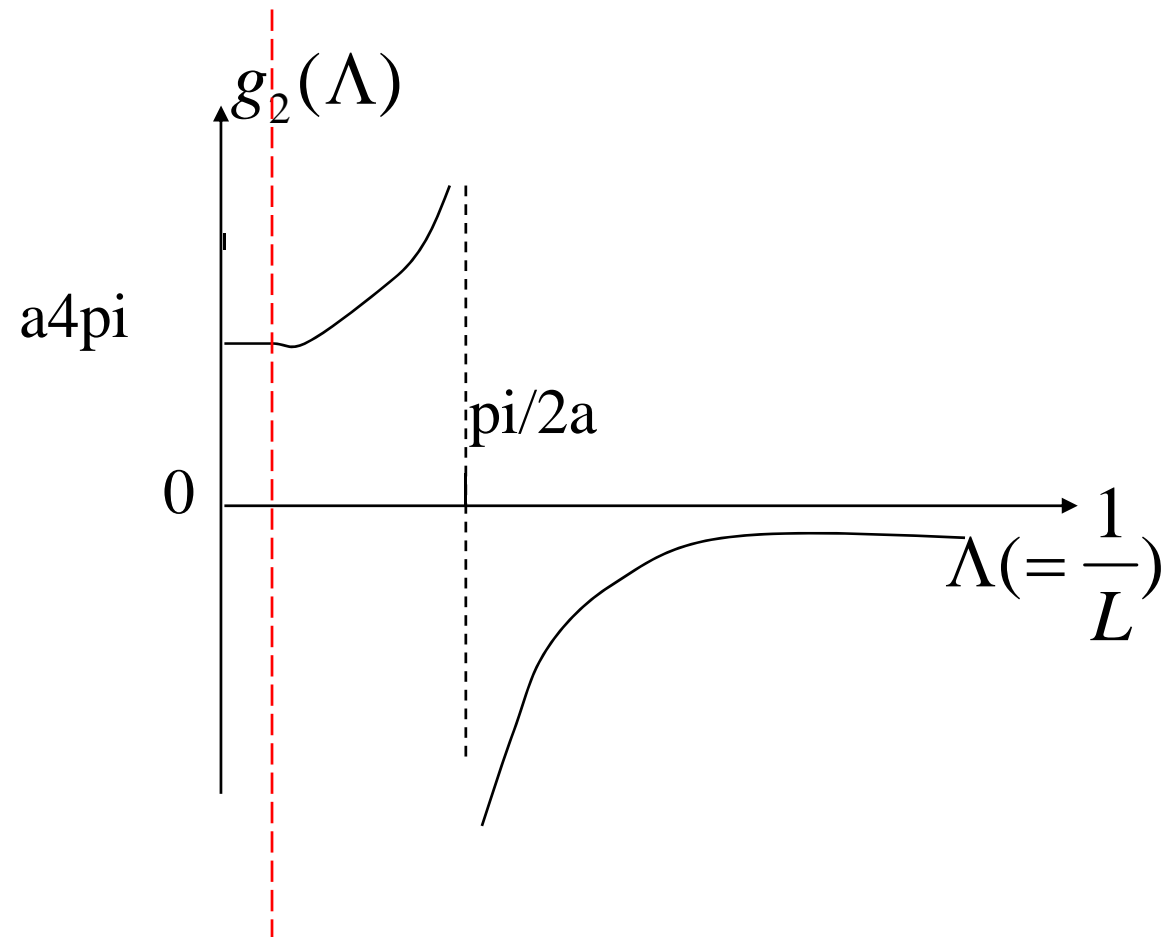
Not equivalent near resonance when $k_F a$ is much larger than unity.

Effective 2-body forces



Borzov, Mashayekhi, Zhang, Song and FZ, to appear 2011

Running 2-body forces

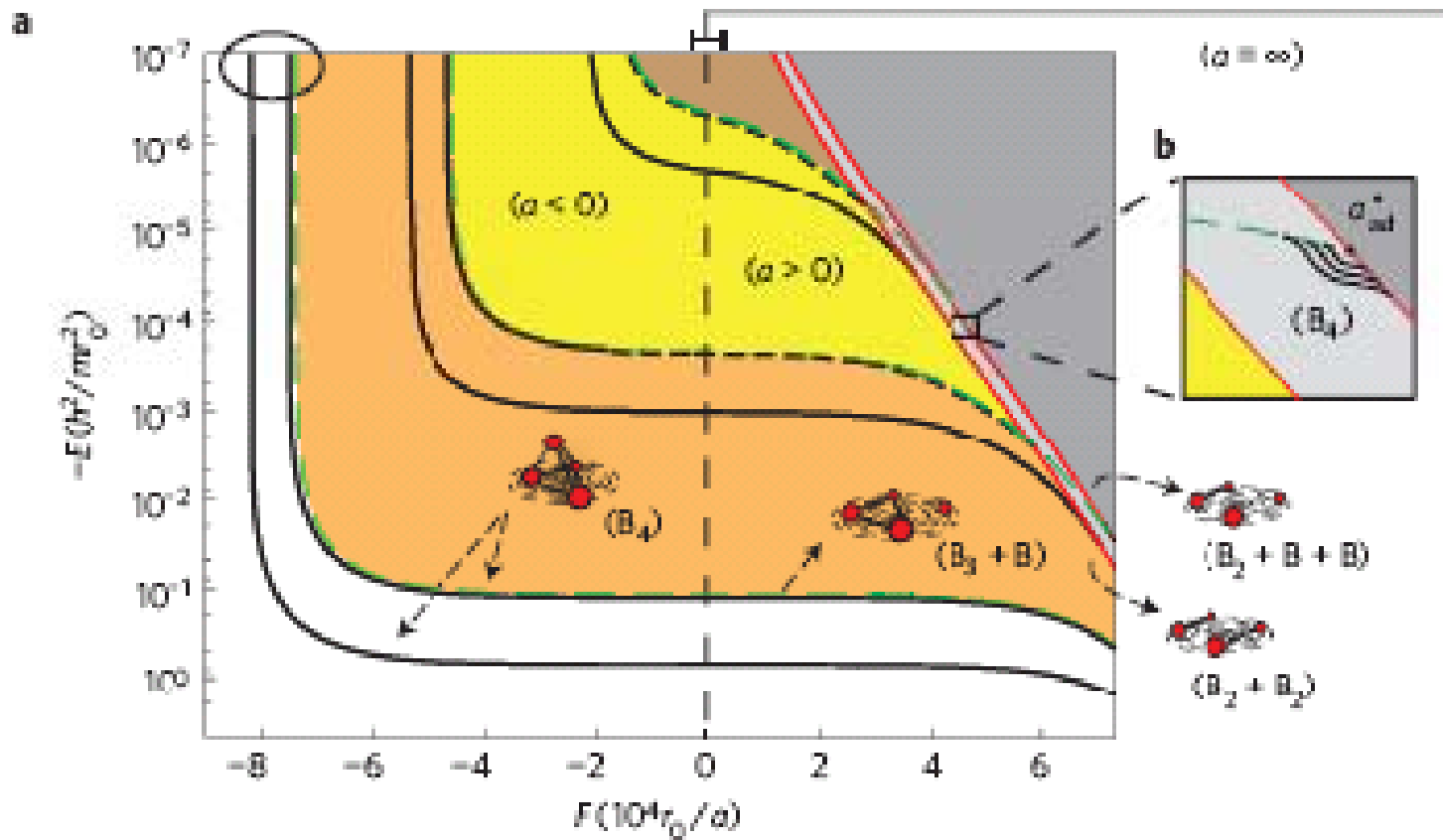


Lecture III

More on RGE

Multiple scatterings in a quantum gas:
Many-body effects

Efimov Trimers (Efimov, 1970-73)



Discovery of Tetramers:

Hammer and Platter Eur.Phys. 2007;

Stecher, D'Incao and C.Greene, Nature Phys 2009.

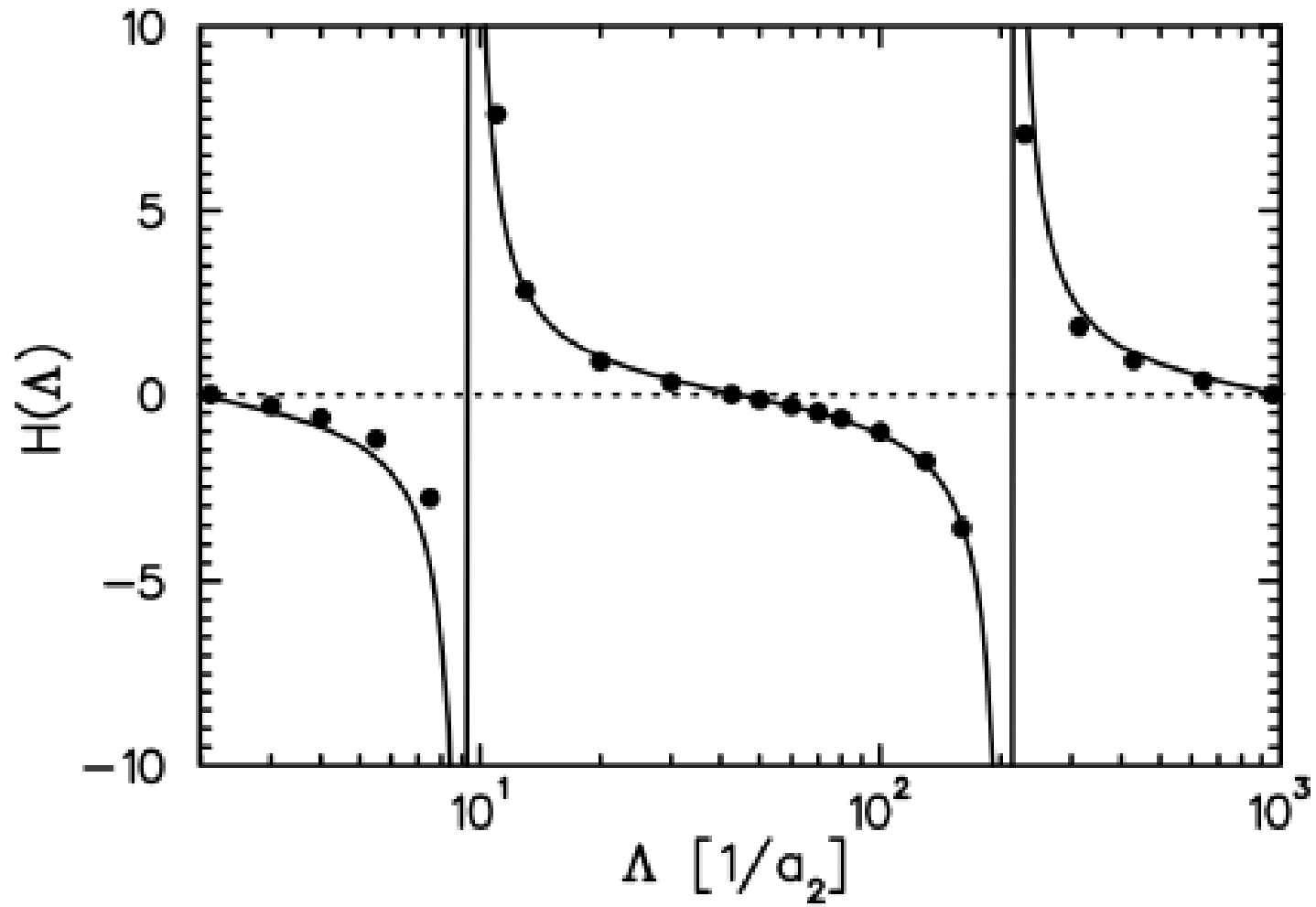
Bedaque-Hammer-van Kolck
three-body Force (Nucl. Phys. A, 98)

$$\begin{aligned} \mathcal{L} = & \psi^\dagger \left(i\partial_0 + \frac{\vec{\nabla}^2}{2M} \right) \psi + \Delta T^\dagger T - \frac{g}{\sqrt{2}} (T^\dagger \psi \psi + \text{h.c.}) \\ & + h T^\dagger T \psi^\dagger \psi + \dots \end{aligned} \quad (2)$$

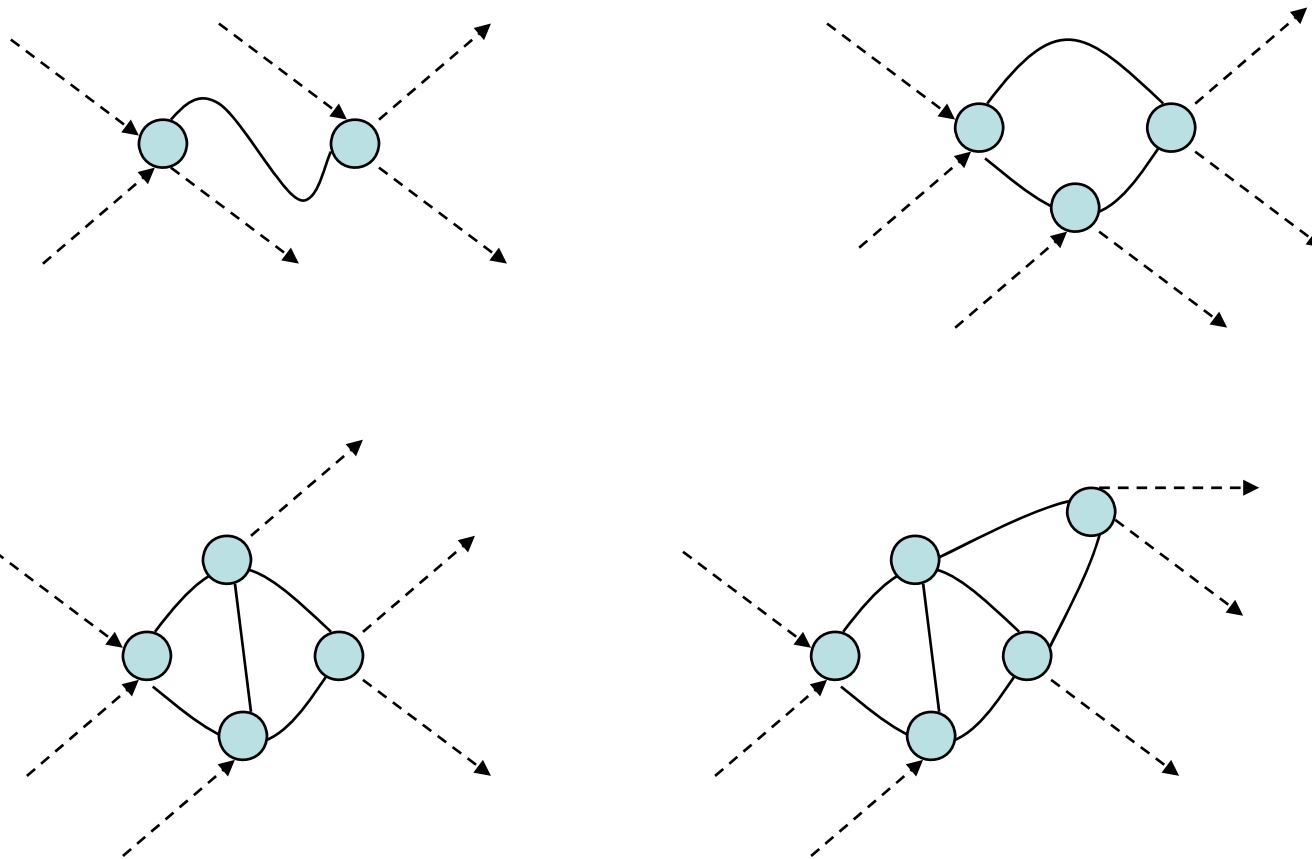
$$H(\Lambda) = - \frac{\sin(s_0 \ln(\frac{\Lambda}{\Lambda_*}) - \text{arctg}(\frac{1}{s_0}))}{\sin(s_0 \ln(\frac{\Lambda}{\Lambda_*}) + \text{arctg}(\frac{1}{s_0}))}.$$

Zero energy atom-dimer scattering length is independent of the cut-off .

Limit cycle RG solution



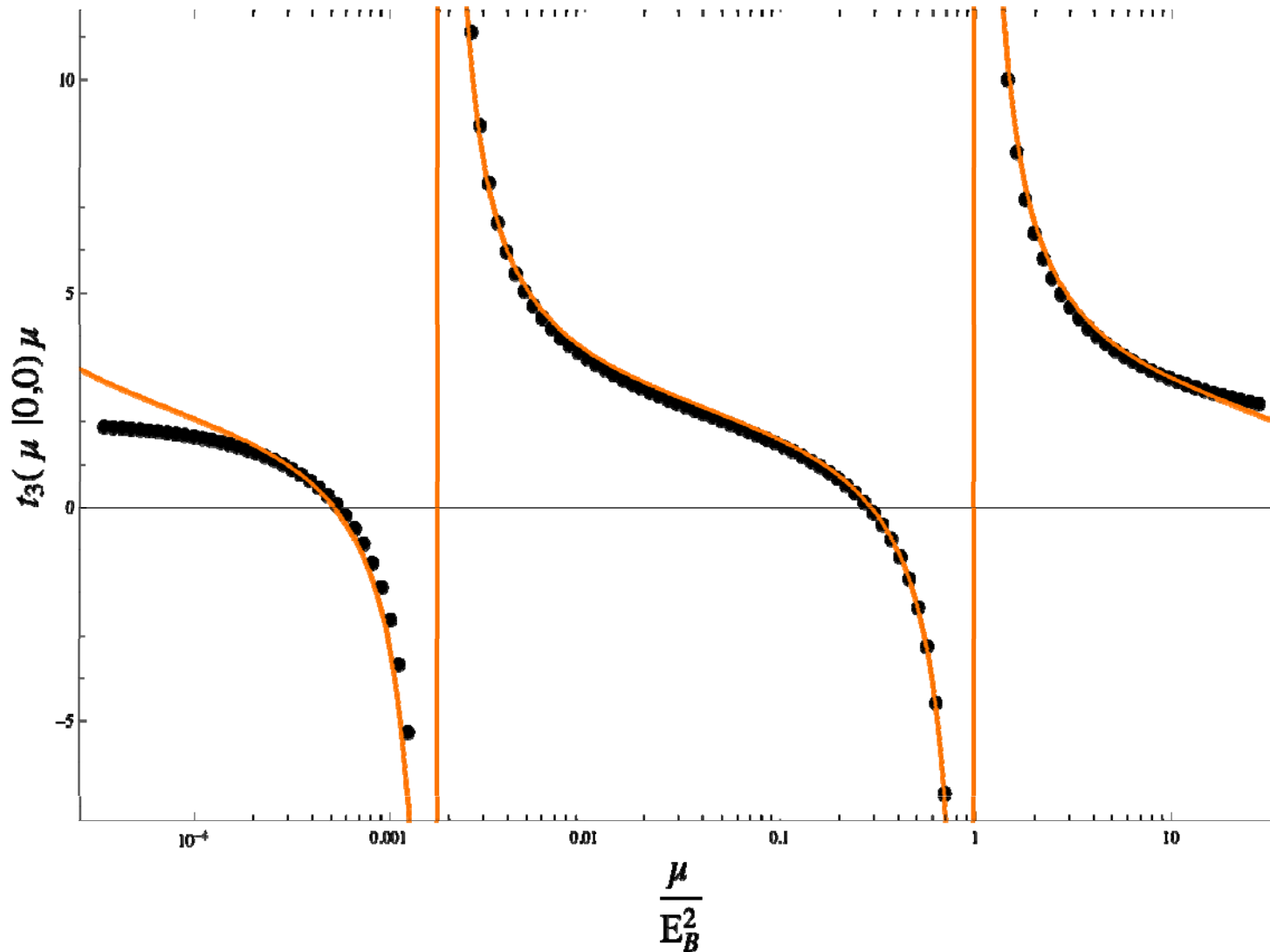
6-pt correlators in BECs: Generalization in a quantum gas



N-loop six-pt correlations (Borzov et al, to appear)

Regularized 6-pt correlator in BECs

$$2.3 + 1.6 \tan\left(4.7 - \frac{\pi}{\log[556]} \log[\mu/E_B^2]\right)$$



A discrete symmetry

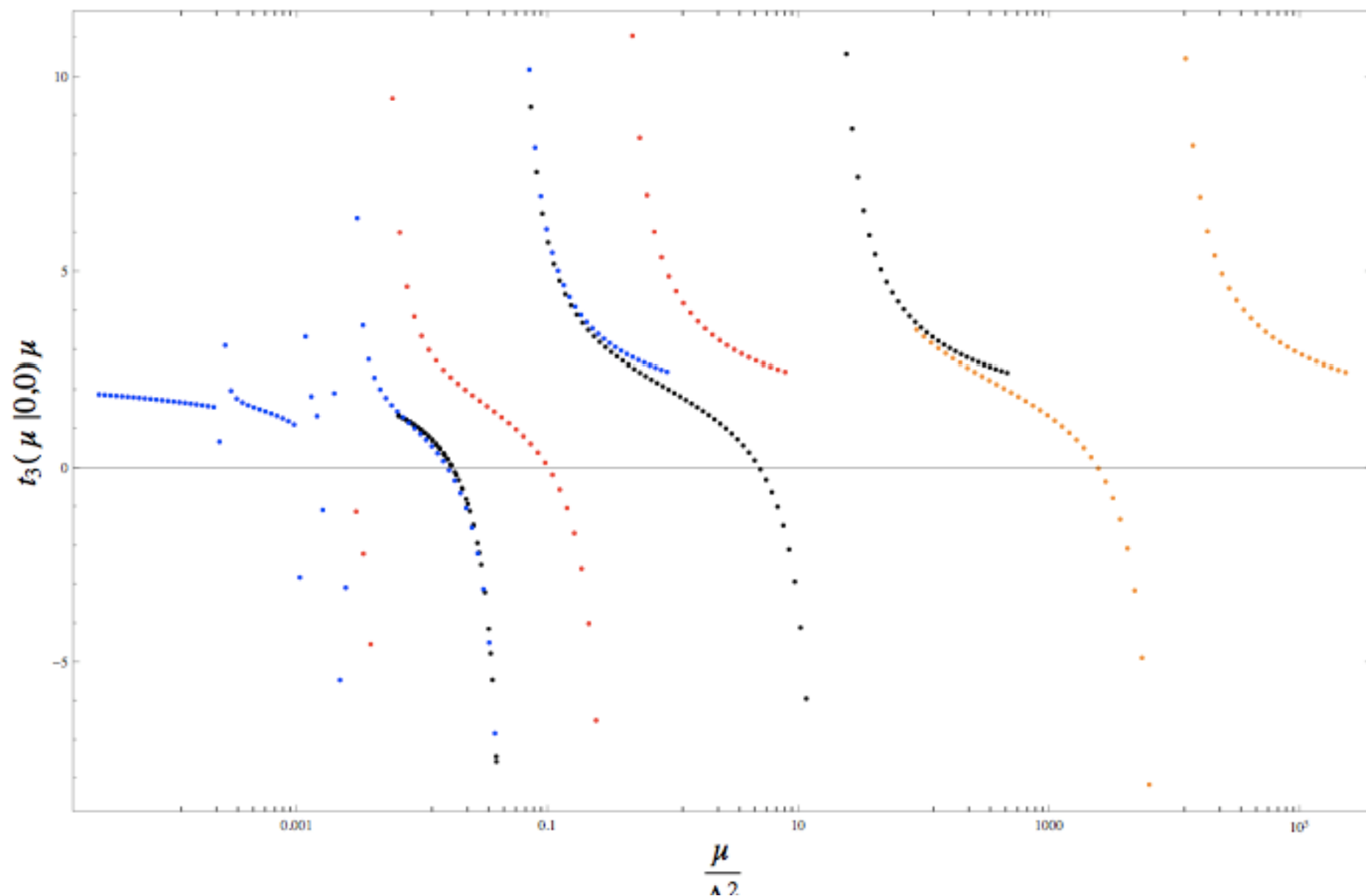
$$G_3(-3\mu\lambda^2; 1/a_2\lambda, \Lambda_3^*) = G_3(-3\mu; 1/a_2, \Lambda_3^*),$$

$$\lambda = e^{\pi/s_0} = 22.7.$$

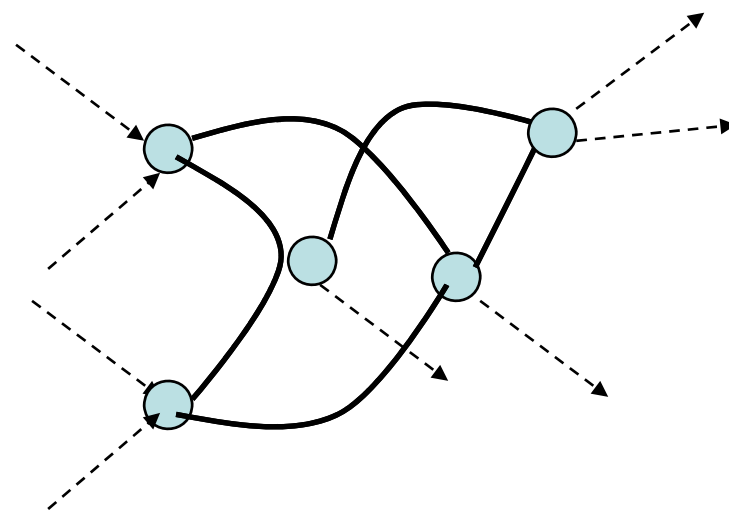
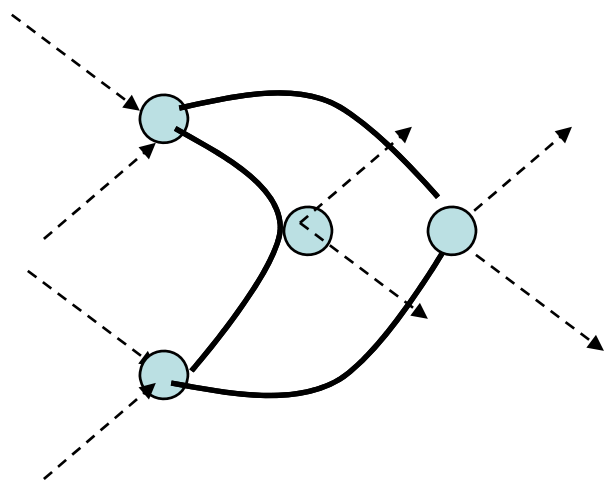
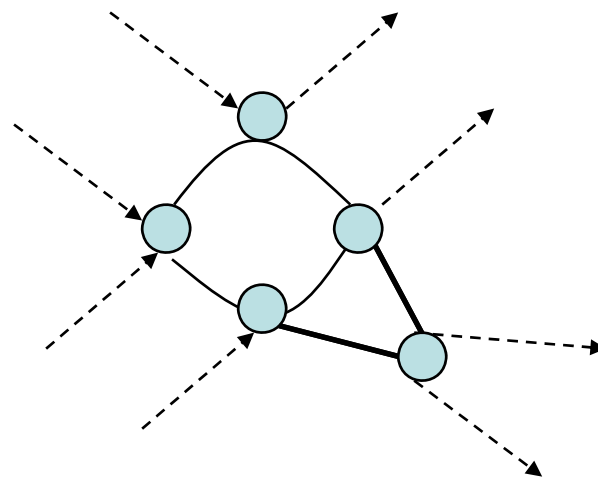
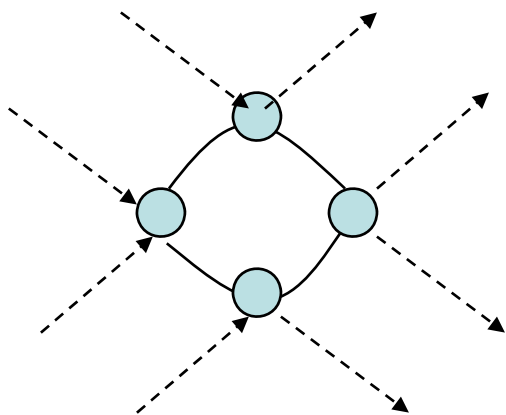
$$G_3(-3\mu\lambda^2; 0, \Lambda_3^*) = G_3(-3\mu; 0, \Lambda_3^*),$$

$$G_3(-3\mu; 0, \lambda^2 \Lambda_3^*) = G_3(-3\mu; 0, \Lambda_3^*),$$

$1/a=0.5$ $t(\mu)$ (Blue) and $1/a=0.5/\lambda$ $t(\lambda^2 \mu)$ (Black) and $1/a=0.5/\lambda^2$ $t(\lambda^4 \mu)$ (Orange) and $1/a=0.5/3$ $t(3^2 \mu)$ (Red, to illustrate that the symmetry is discrete)



4-particle correlators



Issue of multiple scattering in quantum gases

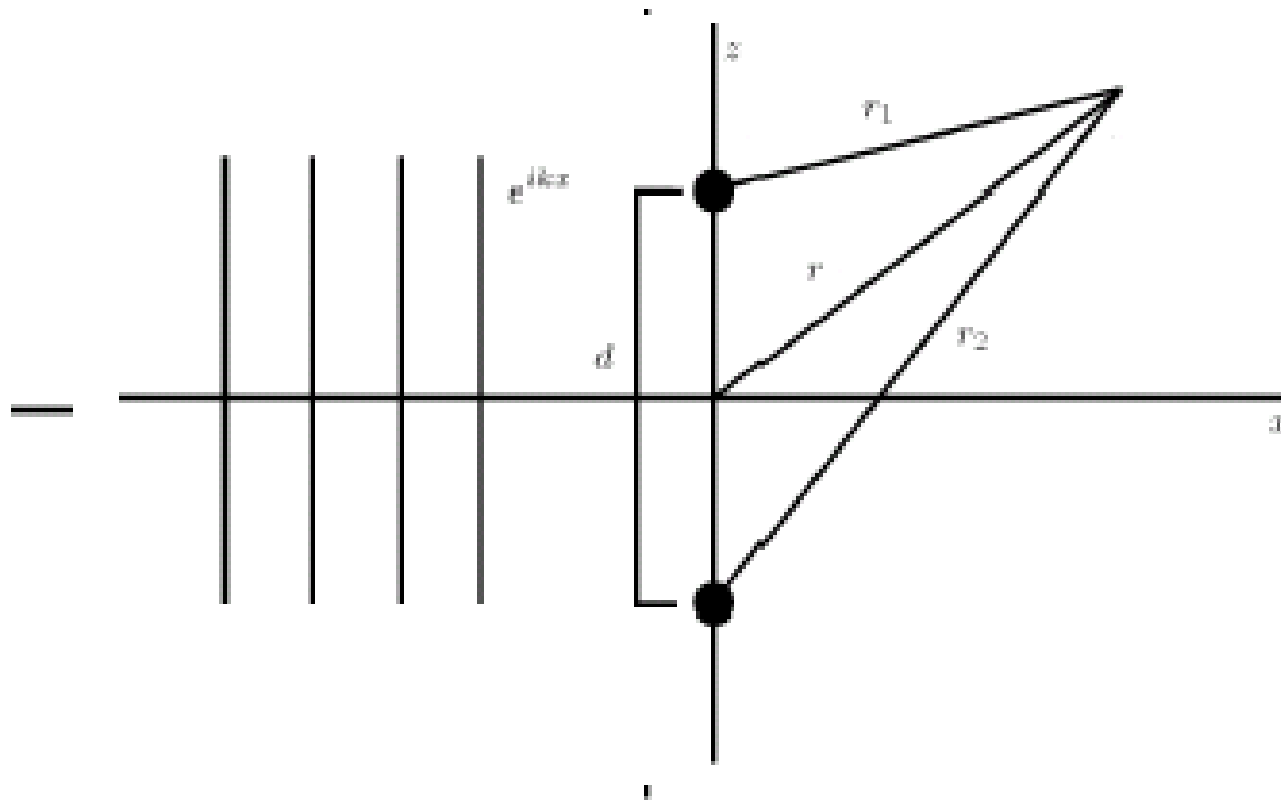
MacNeill and FZ, PRL, 2010

Song and FZ, arXiv. 1104.2924; Phys. Rev. A, 2011;

Song and FZ, 2011, to appear.

- 1) K-space multiple scattering between two atoms when other identical particles are by-standers or as a (statistical) static background.
- 2) Multiple scattering between one of few-body structures and a 3rd particle in quantum gases---dynamical background.

Multiple Scattering ($d \ll a$): From 2 to N ---MacNeill test



Total Cross-section 1)0; 2) 2σ , 3) 4σ , 4) $\sigma \sqrt{\frac{d}{a}}$.

$$\sigma = 4\pi a^2$$

Effective scattering scattering length $\sim d$

“Three-body” bound States

